

Beam divergence near IP and beam-beam effect

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ABP BB-meeting

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Goal of the study

- Evaluate the impact of the betatronic divergence of the 'strong' beam in the presence of a crossing angle and with considering the longitudinal distribution of the bunches (question raised by Stephane in view of HL-LHC)
- The existing 6D lens with crossing-angle disregards the divergence
- An estimator of the importance of the effect is presented

Parameters

	Nominal	High Lum
Bunch population	2×10^{11}	3×10^{11}
β_x^*	0.5 m	0.2 m
β_y^*	0.5 m	0.05 m
σ_z	0.075 m	0.075 m
Φ_{crossing}	160 μrad	720 μrad
ϵ_n	3.75 μm	2.5 μm
σ^*	16 μm	32 μrad
$\sigma_x'^*$	8.2 μm	41 μm

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2D - differential beam-beam kick, gaussian beams

$\sigma_x \neq \sigma_y$ Baseti-Erskine $z1 = x + iy$, $z2 = x \frac{\sigma_y}{\sigma_x} + iy \frac{\sigma_x}{\sigma_y}$

$$dk = \frac{\sqrt{2\pi} N r_0}{A \gamma \sigma_s} \left[w\left(\frac{z1}{B}\right) - e^{-C(x,y)} w\left(\frac{z2}{B}\right) \right] e^{-s^2/2\sigma_s^2} ds$$

$$A = \sqrt{2\pi(\sigma_x^2 - \sigma_y^2)} \quad B = \sqrt{2(\sigma_x^2 - \sigma_y^2)} \quad C(x,y) = \frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2}$$

$w(z)$ is the modified complex error function

Formula valid $\gamma > 0$, $\sigma_x > \sigma_y$, and stable for $(\sigma_x - \sigma_y)/\sqrt{\sigma_x \sigma_y} > 0.01$

$\sigma_x = \sigma_y$ $dk = \sqrt{\frac{\pi}{2}} \frac{N r_0}{\gamma \sigma_s r} \left[1 - e^{-r^2/2\sigma^2} \right] e^{-s^2/2\sigma_s^2} ds$

Small r : $dk = \frac{N r_0 r}{\gamma \sigma^2} \frac{1}{\sqrt{2\pi} \sigma_s} e^{-s^2/2\sigma_s^2} ds$

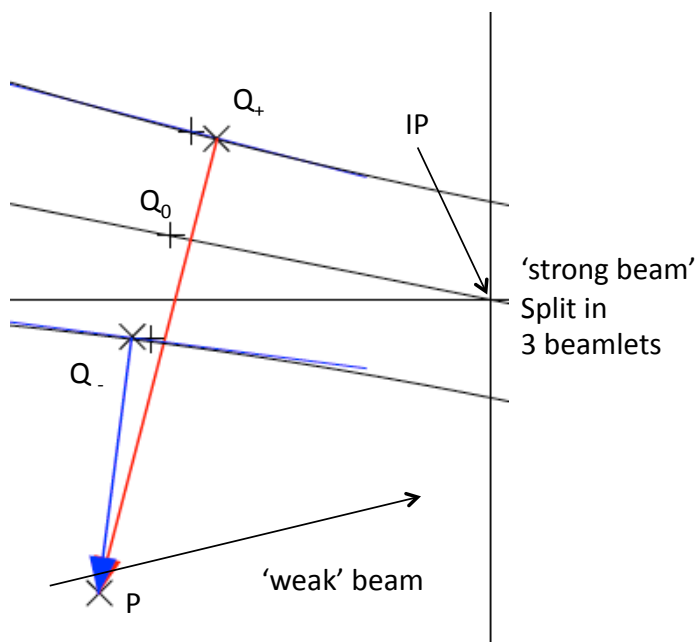
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Method used

- In the crossing plane, the 'Strong' beam is split in 3 beamlets of
 - Null average divergence
 - Negative average divergence
 - Positive average divergence
- Each beamlet is parametrised by its average and r.m.s. value (we can compute a kick properly only for 2D-gaussian beams (Bassetti-Erskine))
- The three corresponding bb-kicks are added
- A tracking is made along the longitudinal coordinate

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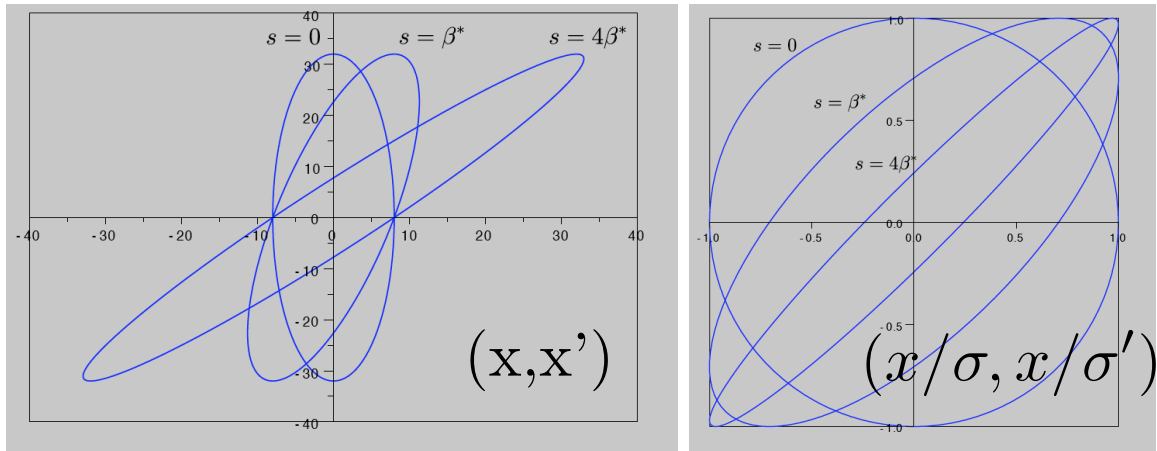
Beam divergence considered or not



- The kick is \perp to the strong beamlet axis
- W/o div, $Q_{0,\pm}$ (+) are aligned with P
- With div (\times) they are not :
 - The distances to P are changed
 - Q_{\pm} move longitudinally, kick intensity is different

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Phase space near the IP, High Lum

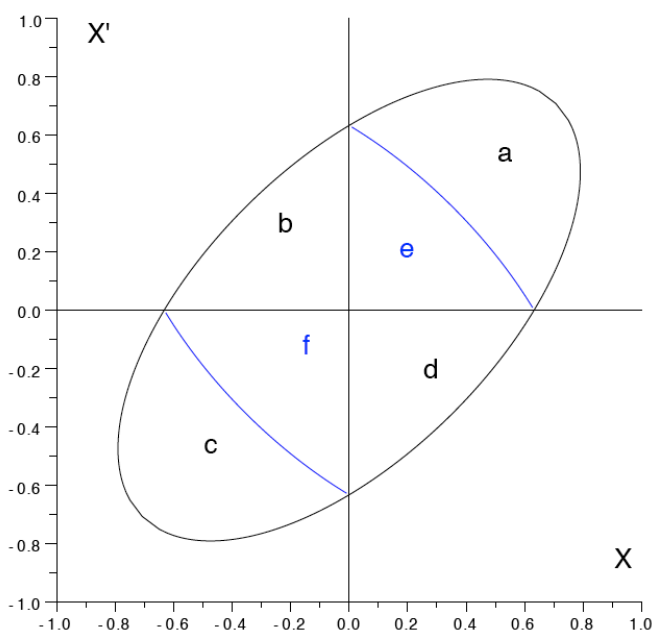


$$\begin{pmatrix} \beta & -\alpha \\ -\alpha & \gamma \end{pmatrix} = \begin{pmatrix} 1 & s \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \beta^* & \alpha^* \\ -\alpha^* & \gamma^* \end{pmatrix} \begin{pmatrix} 1 & 0 \\ s & 1 \end{pmatrix}$$

$$\sigma(s) = \sqrt{\epsilon_n \beta / \gamma_{\text{rel}}} \quad , \quad \sigma'(s) = \sqrt{\epsilon_n \gamma / \gamma_{\text{rel}}}$$

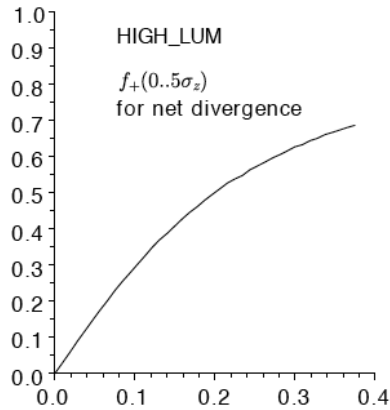
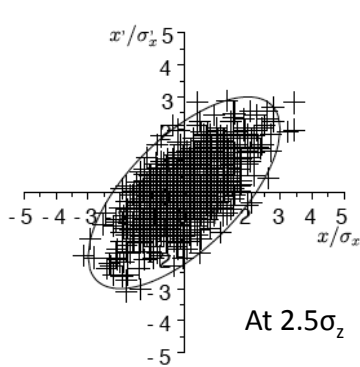
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Building distributions with divergence



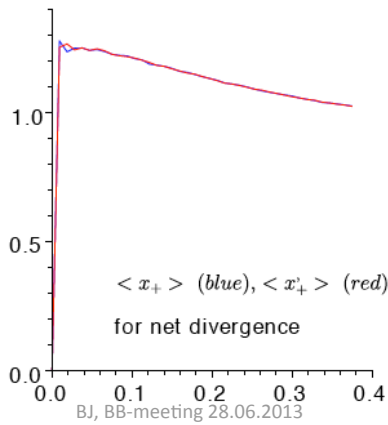
- Build a 2D-gauss distribution with x/x' correlation
- With the 4 sub-samples : a,b,c,d
 - Build e : with d and $x' \rightarrow -x'$
 - Build f : with b and $x \rightarrow -x'$
 - Sub-sample 1 with $\langle x \rangle = \langle x' \rangle = 0$:
central area : b+e+f+d
 - Sub-sample 2 with $\langle x \rangle > 0$ & $\langle x' \rangle > 0$: a-e
 - Sub-sample 3 with $\langle x \rangle < 0$ & $\langle x' \rangle < 0$: c-f
- Get $\langle x \rangle, \langle x' \rangle, \sigma(x), \sigma(x')$ for 1,2,3 as a function of z

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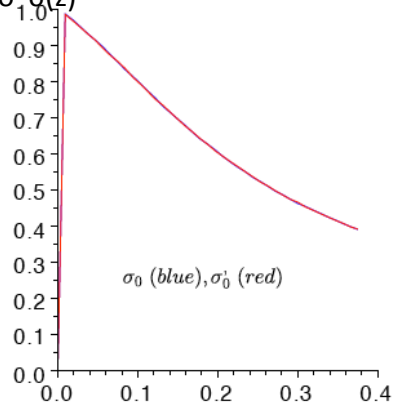
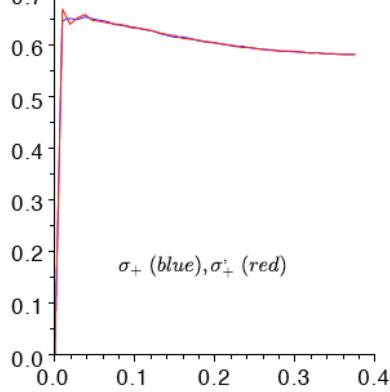


Monte-Carlo filling
of x-x' 'normalized'
phase-space
As a function of z.

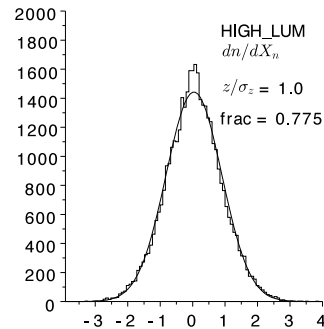
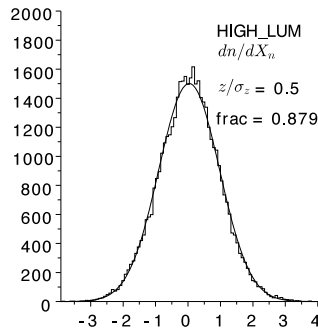
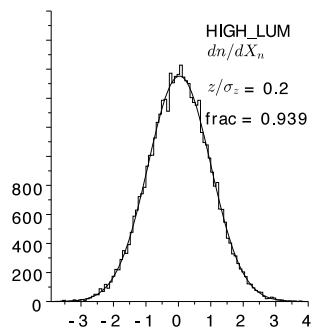
Fraction for
sub-sample 1 : $1-f_+$
sub-sample 2 and 3 : $f_+/2$



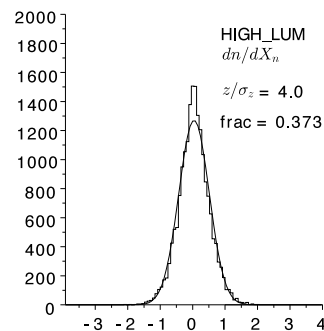
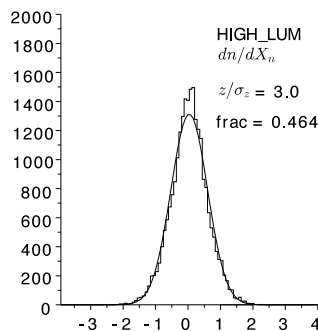
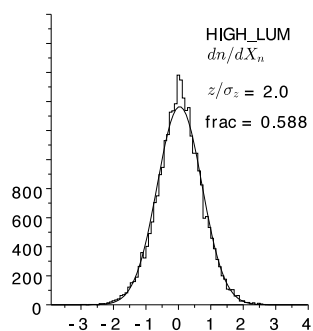
Relative averages and r.m.s. , w.r.t to $\sigma(z)$



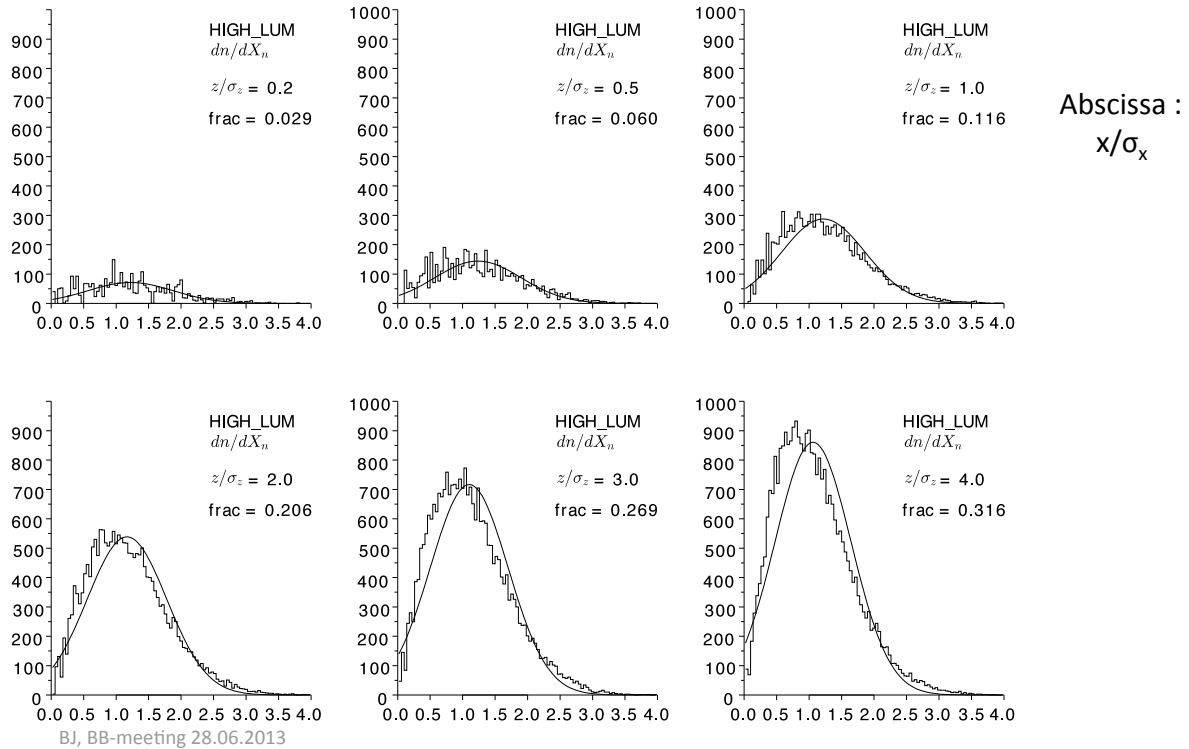
Central beamlet



Abscissa :
 x/σ_x



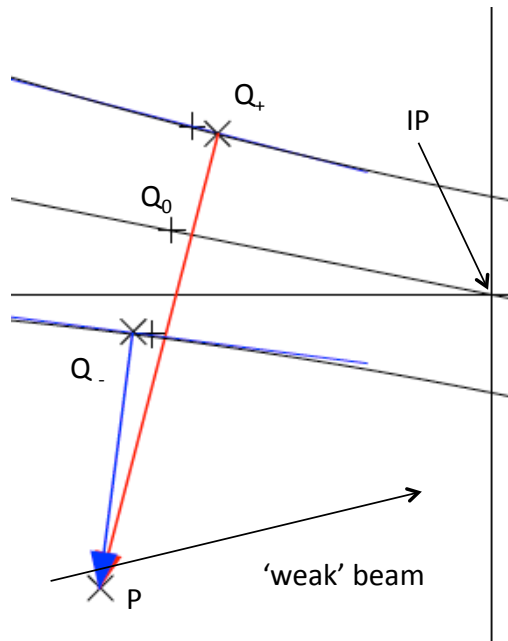
Beamlet of positive divergence



Method used - II

- The decomposition in 3 gaussian beamlets is not perfect
- To compare adequately the two cases (div / no div)
 - The same beamlet decomposition is used for both cases
 - For the no-div case, $\langle x_{\pm}' \rangle = \langle x_0' \rangle = 0$

Tracking



- At P, with $Q_{0,\pm}P$, get

$$d\vec{k} = d\vec{k}_- + d\vec{k}_0 + d\vec{k}_+$$
- Update x' , then x for step ds
- Iterate ...
- This over $-4\sigma_s \rightarrow 4\sigma_s$
- And for $A = [0 \dots 6] \times \sigma_x$ and $\phi = [0 \dots 2\pi]$
- Do everything twice
 - With divergence
 - Without divergence

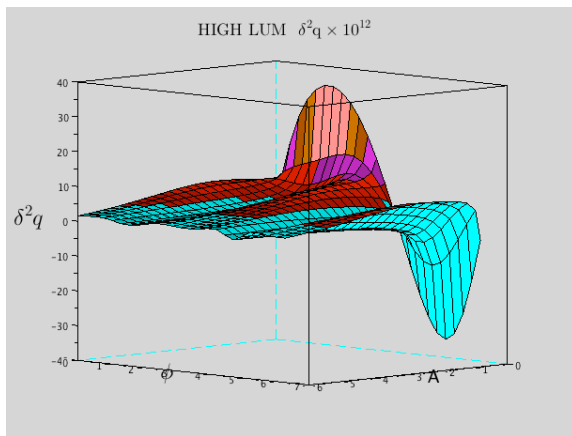
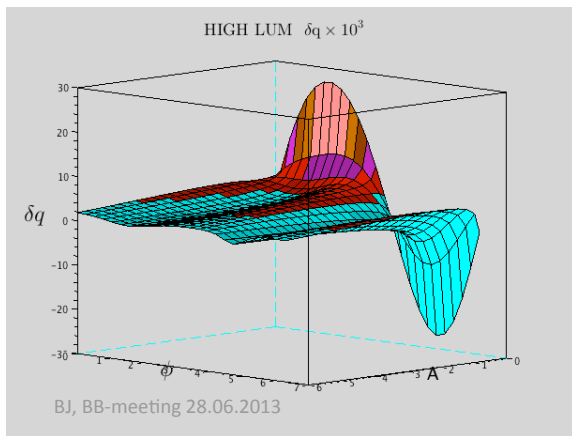
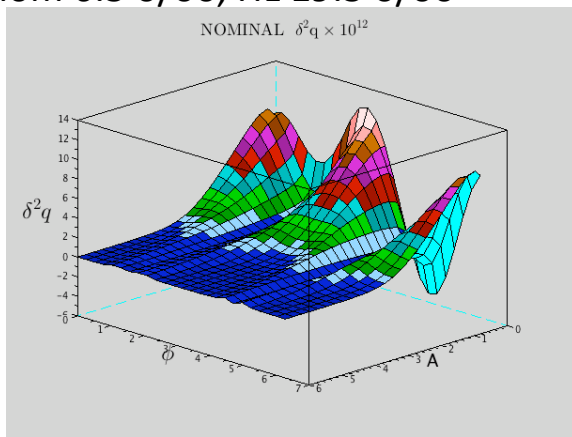
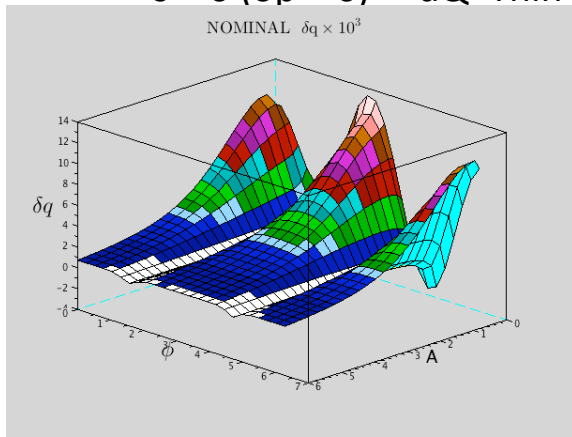
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Results

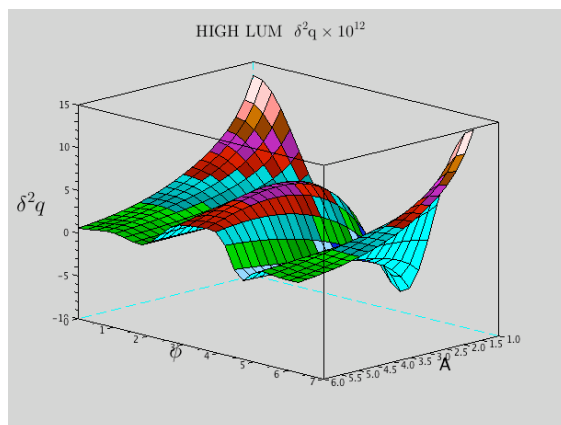
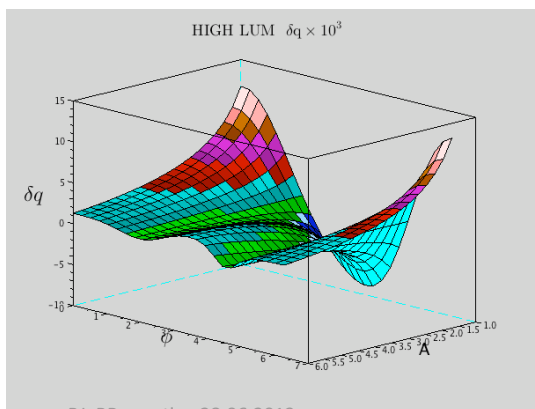
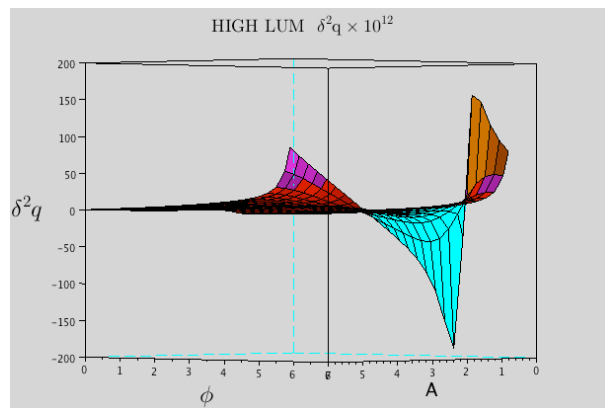
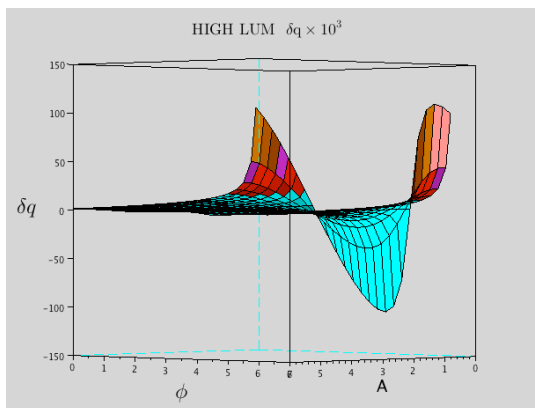
- Start with (x_0, x_0') at IP
- Drift back to $-4\sigma_z$
- Track to $+4\sigma_z$
- Drift back to IP : (x_1, x_1')
- Compute raw δQ as angle between (x_0, x_0') and (x_1, x_1')
- Get $\delta^2 Q = \delta Q_{\text{div}} - \delta Q_{\text{no-div}}$

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$s = 0$ ($\delta p = 0$) dQ_Thin : nom 6.5 o/oo, HL 19.5 o/oo

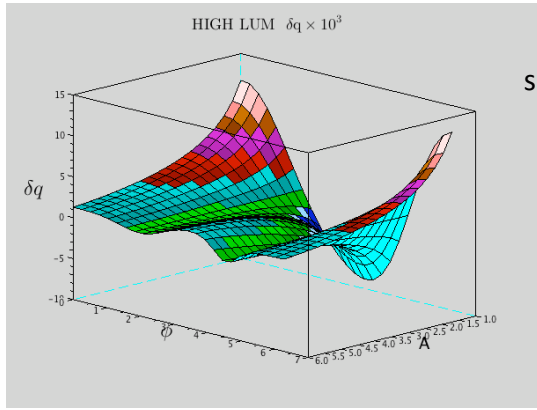


$s = 1\sigma_s$ dQ_Thin : HL 19.5 o/oo

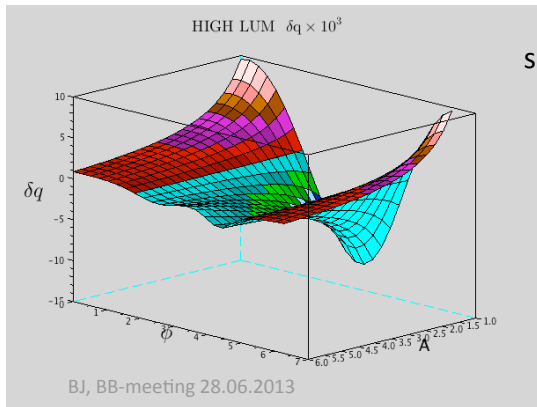
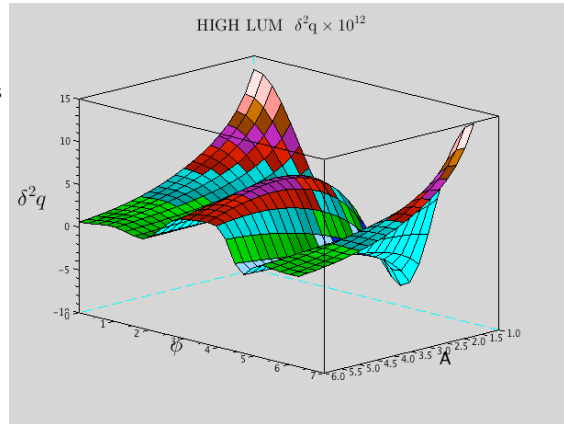


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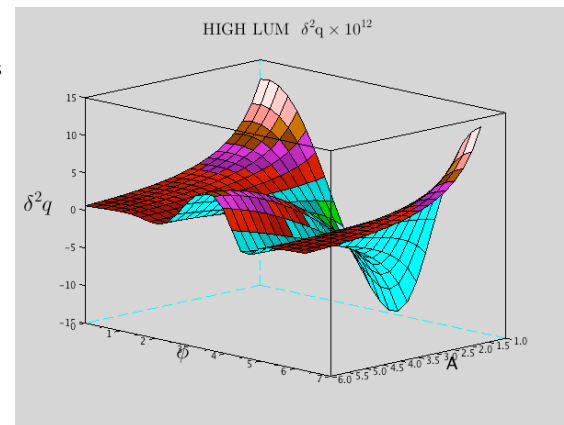
$s = 1,2 \sigma_s \text{ dQ_Thin : HL 19.5 o/o}$



$s = 1\sigma_s$

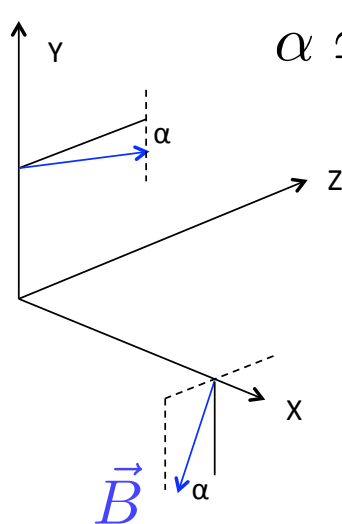


$s = 2\sigma_s$



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Divergence and vertical plane



$$\alpha \simeq \sigma *'_y$$

Diverging fraction, averaged over z : $f_{\text{div}} \simeq 0.5$

$$\delta^2 k = \delta k - \delta k_0 = \frac{f_{\text{div}}}{1 + 2f_{\text{div}}} \frac{\alpha^2}{2} \delta k_0 \simeq \frac{\alpha^2}{8} \delta k_0$$

$$\vec{B} = B_0 \times (0, \cos \alpha, \sin \alpha)$$

\vec{E} Independent of α

$$\Rightarrow \delta k \sim (1 - \alpha^2/4)$$

- Nominal : $\alpha = 3 \times 10^{-5} \rightarrow \alpha^2/8 = 0.13 \times 10^{-9}$
- High-Lum : $\alpha = 8 \times 10^{-5} \rightarrow \alpha^2/8 = 0.80 \times 10^{-9}$

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Results about divergence

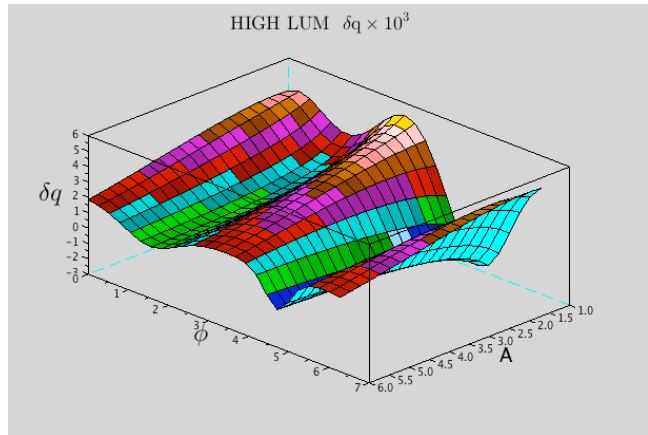
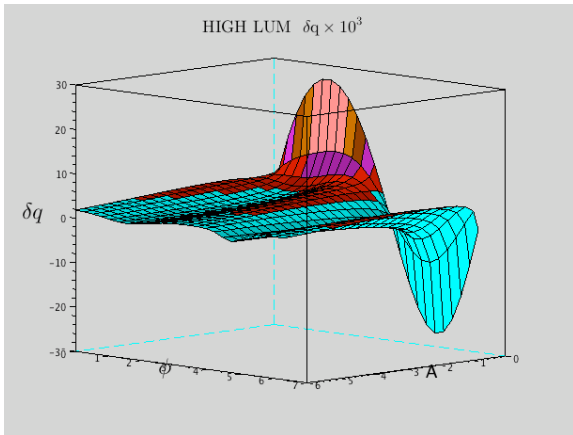
- Considering the tune variations
 - In the crossing plane :
 - The difference between the two cases, divergence considered or not considered is $\delta^2 q_{\text{rel}} < 2 \times 10^{-9}$, both with NOMINAL and HL.
 - This difference similar when the average longitudinal position of the test particle w.r.t. to the strong bunch is changed ($0 \rightarrow 2\sigma_s$).
 - In the other plane :
 - The effect is 10× smaller with NOMINAL ,i.e. $\delta^2 q_{\text{rel}} \approx 0.13 \times 10^{-9}$
 - The effect is 2× smaller with HL ,i.e. $\delta^2 q_{\text{rel}} \approx 0.8 \times 10^{-9}$

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Small amplitude distortions,
(independent of divergence effect)

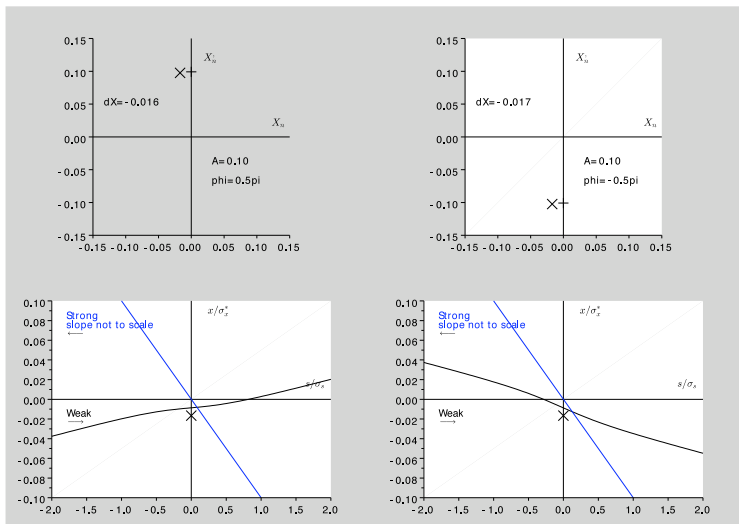
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The apparent dQ excursion at small amplitude and phase space angle $\pm\pi/2$



- The raw dQ grows without limit at towards small amplitude ($x=0, x' \neq 0$), particularly marked with HL
- What happens ?

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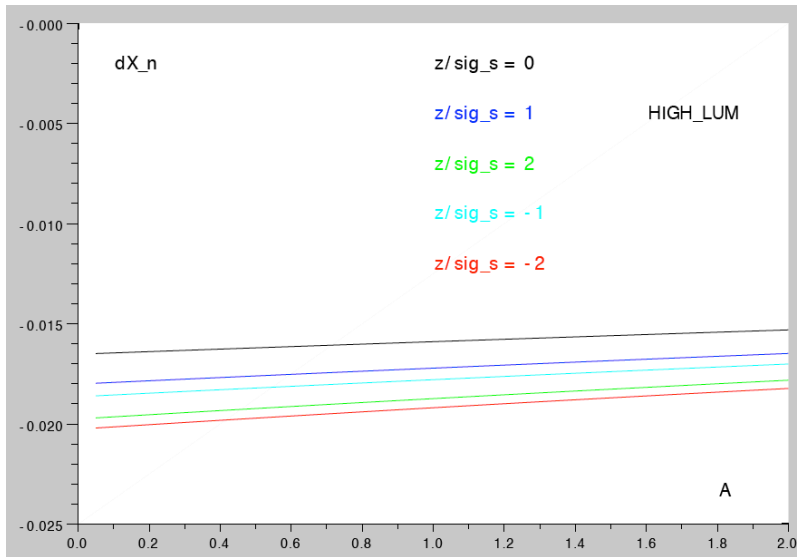


High_LUM

- With bunch length considered and crossing angle
 - A δX appear with tacking, with the same sign whatever the phase angle
 - So, this is an orbit effect
 - At High-LUM : $\delta X = -0.016 \sigma^*$
 - The same applies to the other beam \rightarrow collision mismatch of $3\% \sigma^*$
 - Problematic ?

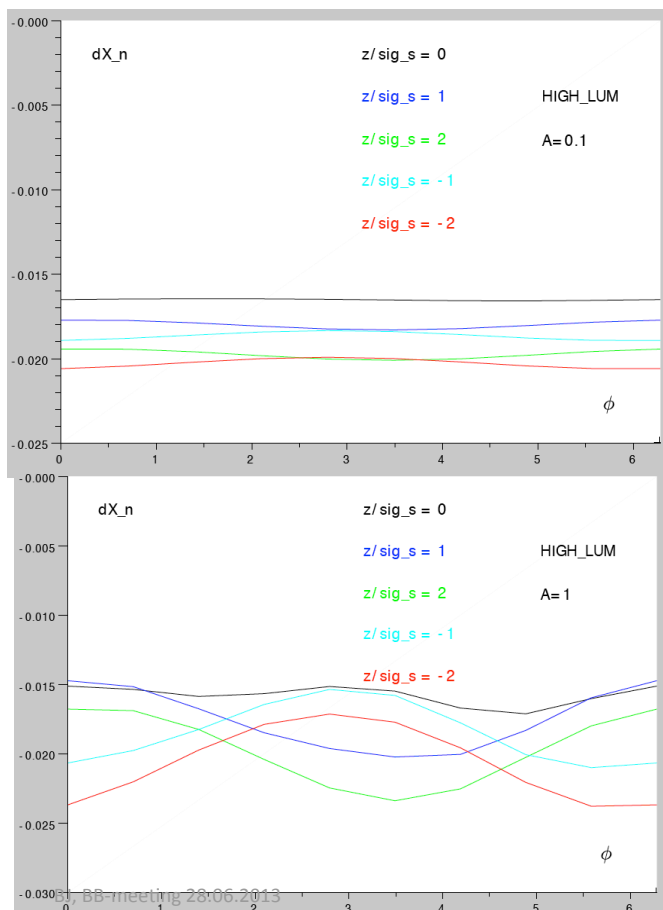
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Beam displacement High-LUM




- Vary A , $\Phi = \pi/2$
- Slight variations
 - with amplitude
 - and z -displacement (δp)

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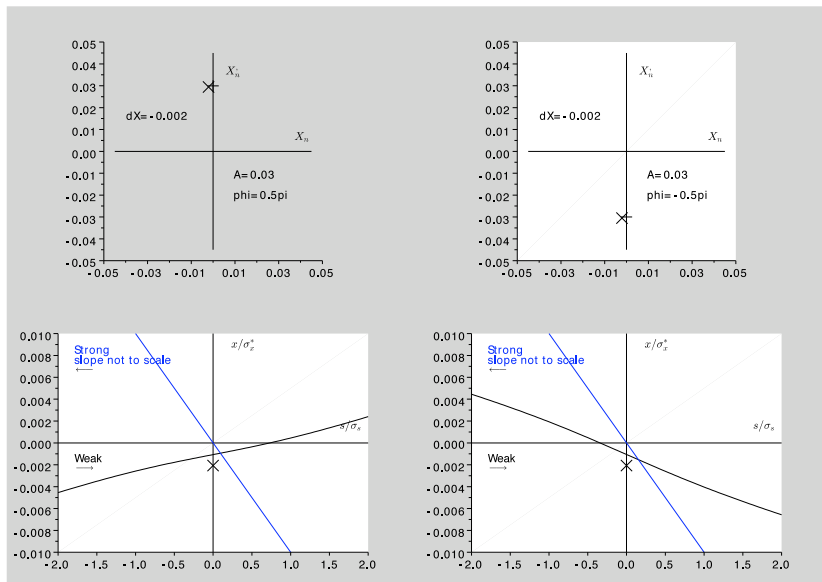
Beam displacement High-LUM - II

- Vary Φ
- Slight variations
 - with Φ (1% of σ)
 - with z -displacement (δp)


 (2 ± 0.5) % σ^* at $A=1$

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Nominal



- Much smaller effect, of $2 \text{ o}/\infty \sigma^*$

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Summary

- The effect of the betatronic divergence can be safely neglected ($\delta^2 Q / \delta Q < 2 \times 10^{-9}$) with both nominal and high luminosity collision parameters
- Small orbit effect (partly amplitude & phase dependent) visible with 'thick lens' beam-beam tracking

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