



Beta Beams, EUROnu WP4



One of the Beta Beam Challenges:

Collective Effects



Christian Hansen

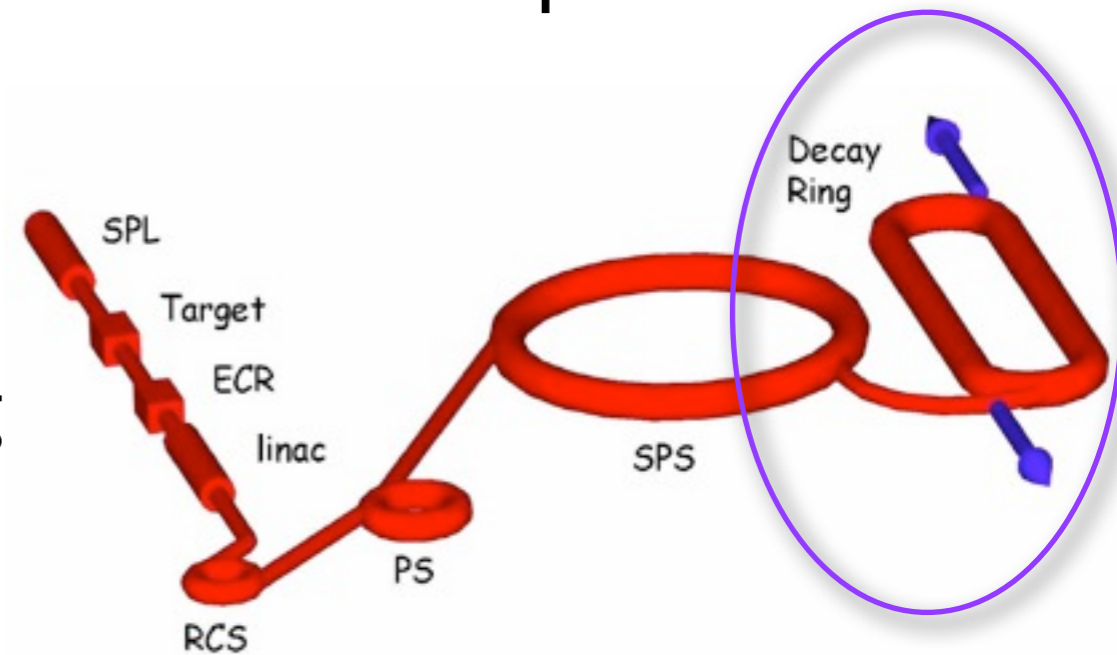
20/01/2011, EUROnu Annual Meeting, RAL

Many thanks to: E. Benedetto, A. Chancé, K. Li, E. Metral, N. Mounet, G. Rumolo, B. Salvant & E. Wildner

BB Collective Effects Studies

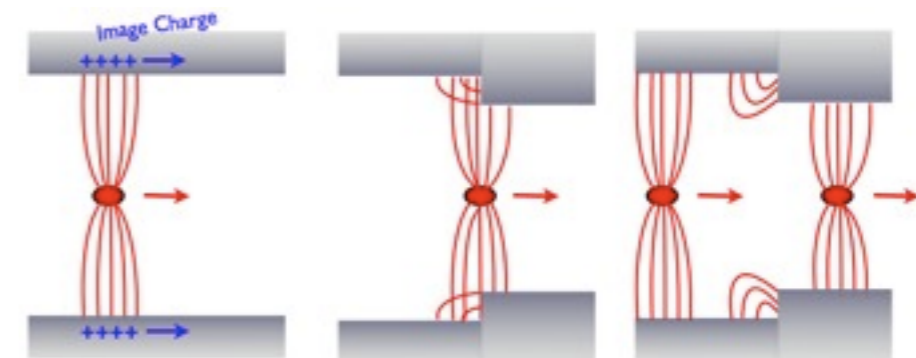
- Instability studies are important for the Beta Beam project, since
 - ➔ High intensity ion beams are foreseen
 - ➔ Collective Effects could limit the final performance

- Will study all machines
 - ➔ Today only the Decay Ring



- Will study all possible reasons for instabilities
 - ➔ Today only results from single bunch

“Transverse Resonance Broad Band Impedance”



3 Tools

- Three ways to find the Bunch Intensity Limit, N_b^{th} :
 - ➔ A multi-particle tracking program in time domain, “HEADTAIL”
 - ➔ A theoretical program in frequency domain, “MOSES”
 - ➔ Peak current values into a coasting beam formula gives the “Coasting Beam Equation” (here for $\xi=0$):

G. Rumolo et al,
CERN-SL-
Note2002-036-AP

Y.H.Chin CERN-
LEP-TH/88-05

$$N_{b_{x,y}}^{th} = \frac{32}{3\sqrt{2}\pi} \frac{R|\eta|\varepsilon_l^{2\sigma}\omega_r}{\langle\beta\rangle_{x,y} Z^2\beta^2 c R_{\perp}}$$

E. Métral, CERN,
Overview of Single-Beam
Coherent Instabilities in
Circular Accelerators

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E. Métral, CERN,
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R_{\perp} = “Shunt Impedance” (see next slide)

R_{\perp} of the DR

Private Discussions
with G. Rumolo

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- Detailed calculations of Transversal Shunt Impedance, R_{\perp} , require design assumptions of ALL DR components, instead:
- **Let's estimate R_{\perp}^{DR} based on a machine with same circumference as DR; SPS ($R_{\perp}^{SPS} = 20 \text{ M}\Omega/\text{m}$)**
 - ➔ **Modern, smooth design of the vacuum pipe compare to old SPS → *Improvement by factor 10***
 - $R_{\perp}^{DR} \sim 2 \text{ M}\Omega/\text{m}$
 - ➔ **The DR is a less general machine than the SPS (not required to handle many type of beams)**
 - ➔ **No need for as many kickers as SPS (and modern kicker design) → *Improvement by factor 2***
 - $R_{\perp}^{DR} \sim 1 \text{ M}\Omega/\text{m}$
- Further; in 20 years improved Broad Band Feedback System

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N_b^{th} vs. R_{\perp} in DR

C. Hansen, CERN, Nufact10
Collective Effect Studies of a
Beta Beam Decay Ring

- **According to the Coasting Beam Equation (CB Eq.) R_{\perp} is the only parameter not fixed by FP6 design**
- **Let's find required shunt impedance, R_{\perp}^{req} ;**

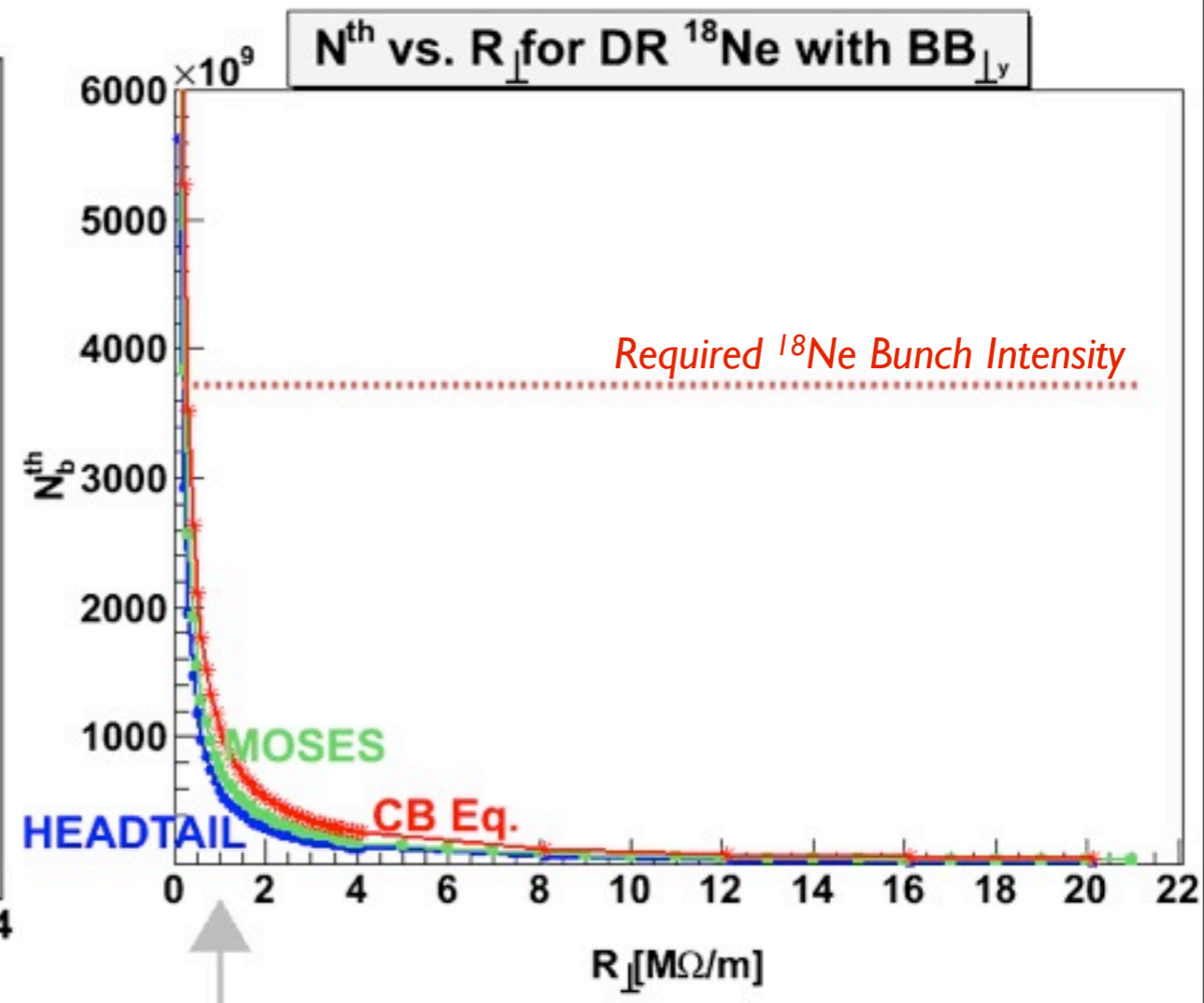
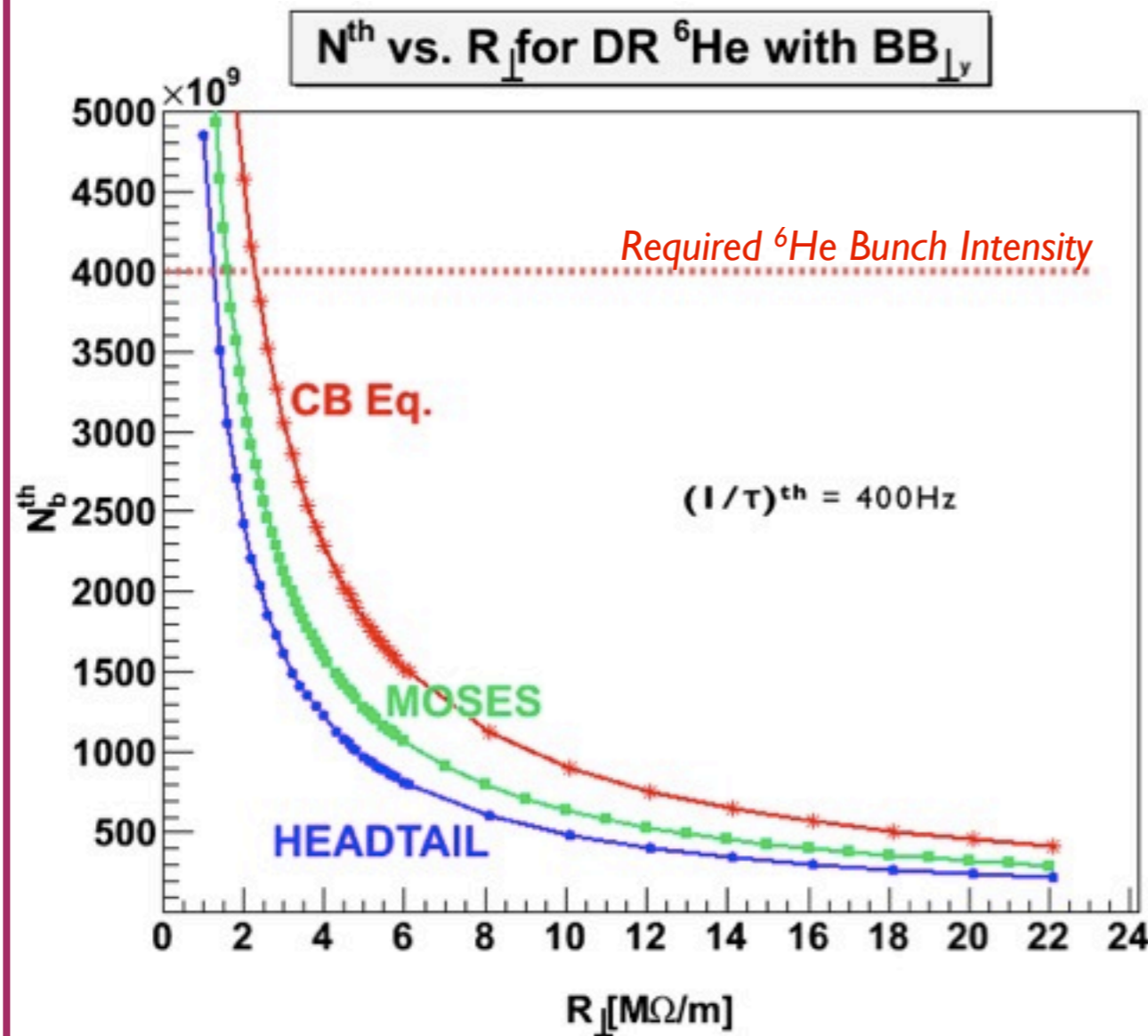
T H R E S H O L D

N_b^{th} vs. R_{\perp} in DR

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THRESHOLD

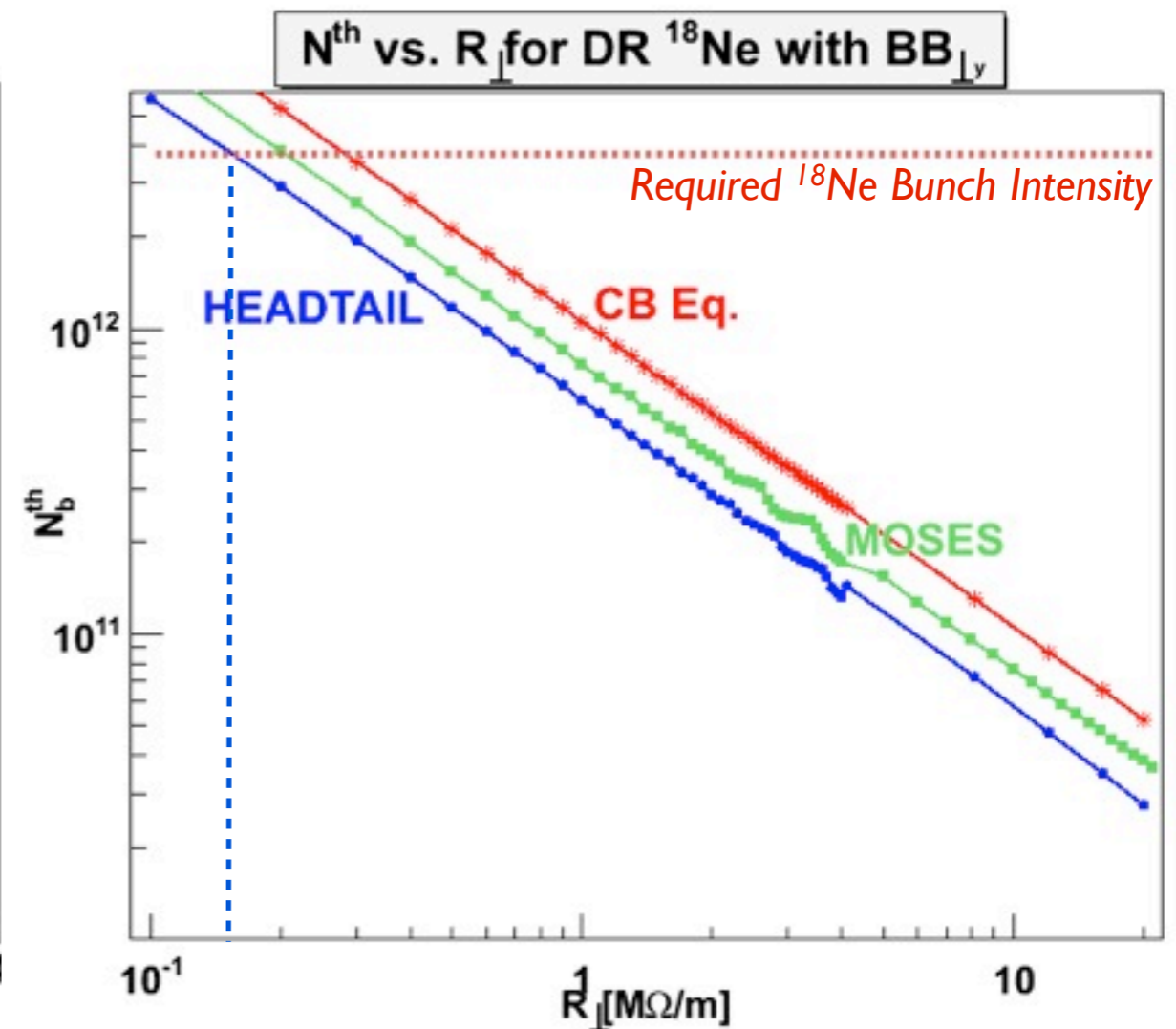
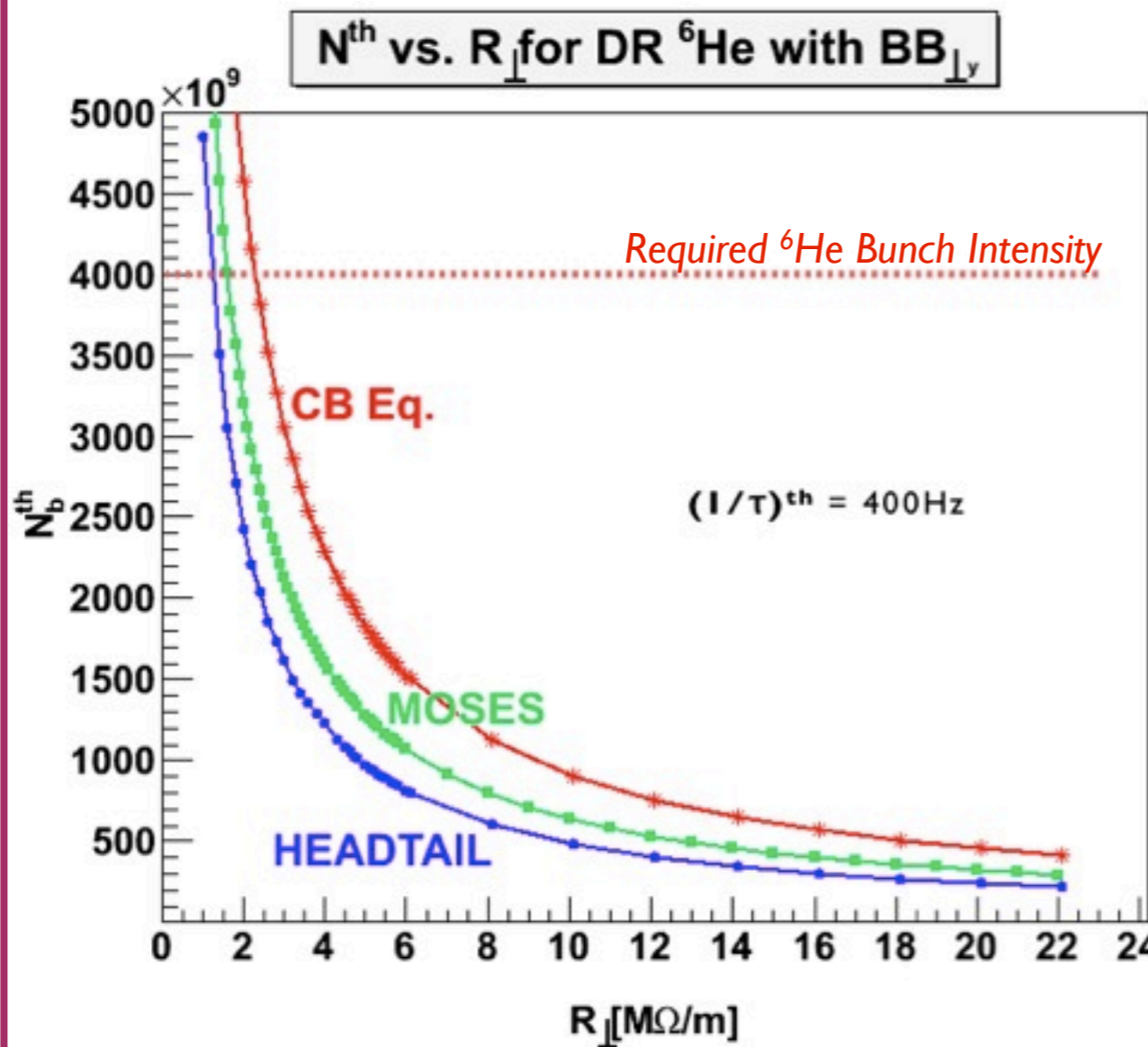


N_b^{th} vs. R_{\perp} in DR

C. Hansen, CERN, Nufact10
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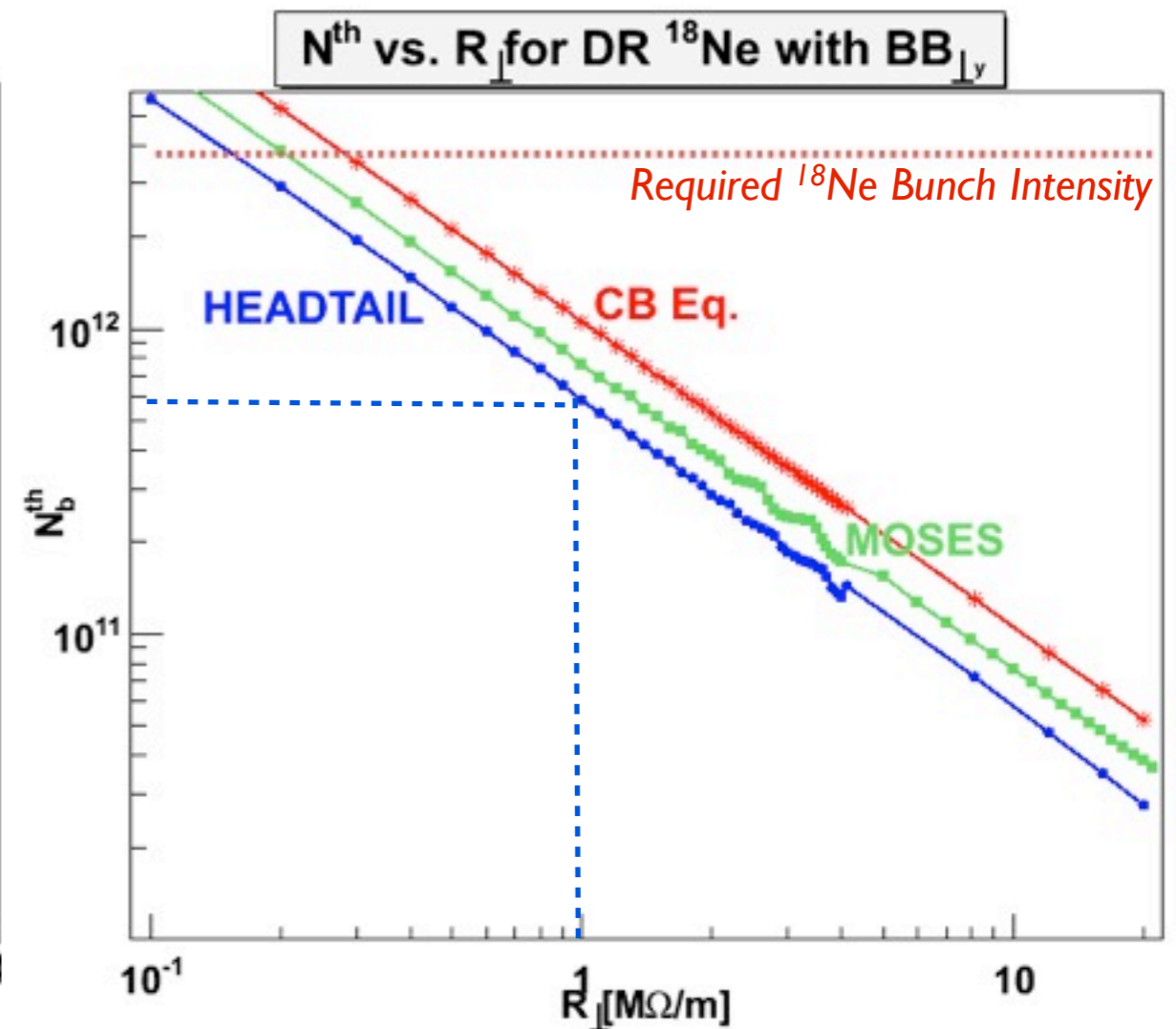
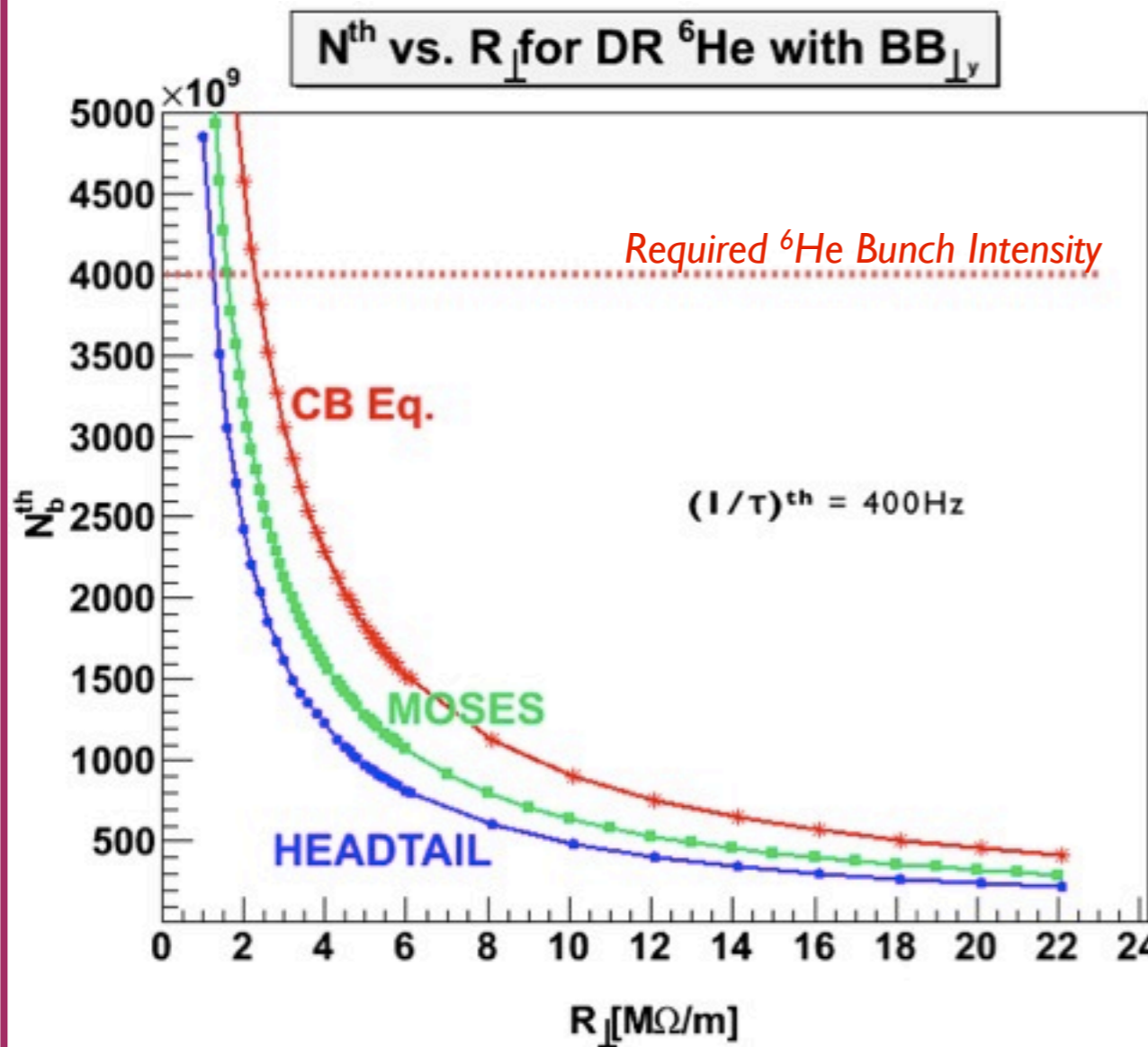
- $R_{\perp}^{req} = 0.15 \text{ M}\Omega/\text{m}$ to allow $N_b^{th} = 3 \times 10^{12} \text{ }^{18}\text{Ne}$
So since $R_{\perp}^{req} < R_{\perp}^{DR} \sim 1 \text{ M}\Omega/\text{m} \rightarrow$ Redesign of DR

N_b^{th} vs. R_{\perp} in DR

C. Hansen, CERN, Nufact10
Collective Effect Studies of a
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- According to the Coasting Beam Equation (CB Eq.) R_{\perp} is the only parameter not fixed by FP6 design
- Let's find required shunt impedance, R_{\perp}^{req} ;

THRESHOLD



- $N_b^{th} = 6e11$ ${}^{18}\text{Ne}$ can be used to get N_b^{th} for all other ions by using that CB Eq. goes as

$$N_{b_{x,y}}^{th} \propto \frac{A}{Z^2}$$

DR Intensity Limit for FP6 Lattice

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$R_{\perp DR} =$ 1 MΩ/m	Bunch Intensity Limit, N_b^{th}	
	[e12]	[% of N_b^{nom}]
${}^6\text{He}$	5.0	100
${}^{18}\text{Ne}$	0.6	16
${}^6\text{He}$	5.0	52
${}^{18}\text{Ne}$	0.6	32
${}^6\text{He}$	5.0	52
${}^{18}\text{Ne}$	0.6	81
${}^8\text{Li}$	3.0	60
${}^8\text{B}$	1.1	59
${}^8\text{Li}$	3.0	30
${}^8\text{B}$	1.1	29
${}^8\text{Li}$	3.0	12
${}^8\text{B}$	1.1	12

Note; In Donini's table $SF = 10^{-4}$ while we are using $SF = 5 \cdot 10^{-3}$

A. Donini, Summary on Beta-Beams

Ions	Fluxes [10^{18}]	Years	$(\sin^2 2\theta_{13})_{min}$	NH, $(\sin^2 2\theta_{13})_{min}$
${}^6\text{He}$	$\bar{\Phi}_0 = 2.9$	5	5×10^{-4}	No Sensitivity
${}^{18}\text{Ne}$	$\Phi_0 = 1.1$	5		
${}^6\text{He}$	$\bar{\Phi}_0 \times 2$	2	6×10^{-4}	No Sensitivity
${}^{18}\text{Ne}$	$\Phi_0/2$	8		
${}^6\text{He}$	$\bar{\Phi}_0 \times 2$	2	1×10^{-3}	No Sensitivity
${}^{18}\text{Ne}$	$\Phi_0/5$	8		
${}^8\text{Li}$	$\bar{\Phi}_0$	5	1.5×10^{-3}	3×10^{-2}
${}^8\text{B}$	Φ_0	5		
${}^8\text{Li}$	$\bar{\Phi}_0 \times 2$	5	7×10^{-4}	1.5×10^{-2}
${}^8\text{B}$	$\Phi_0 \times 2$	5		
${}^8\text{Li}$	$\bar{\Phi}_0 \times 5$	5	2×10^{-4}	8×10^{-3}
${}^8\text{B}$	$\Phi_0 \times 5$	5		
${}^8\text{Li}$	$\bar{\Phi}_0$	3	1.7×10^{-3}	3×10^{-2}
${}^8\text{B}$	Φ_0	5		
${}^6\text{He}$	$\bar{\Phi}_0$	2		
${}^8\text{Li}$	$\bar{\Phi}_0 \times 2$	3	7×10^{-4}	1.5×10^{-2}
${}^8\text{B}$	$\Phi_0 \times 2$	5		
${}^6\text{He}$	$\bar{\Phi}_0 \times 2$	2		
${}^8\text{Li}$	$\bar{\Phi}_0 \times 5$	3	3×10^{-4}	8×10^{-3}
${}^8\text{B}$	$\Phi_0 \times 5$	5		
${}^6\text{He}$	$\bar{\Phi}_0 \times 5$	2		

${}^8\text{Li}$	${}^8\text{B}$	${}^6\text{He}$
60	59	100
30	29	50
12	12	21

$$N_b^{nom} = \frac{\Phi_0 L_{circ} t_{sps}}{N_{bunches} L_{eff} T_{eff}} \left(1 - 2^{-\frac{t_{sps}}{\gamma t_{1/2}}}\right)^{-1}$$



Beyond FP6

Decay Ring Redesign

DR REDESIGN

- So far all studies based on **EURISOL** FP6 parameters

- According to CB Eq.

Data Base:
<http://j2eeps.cern.ch/beta-beam-parameters/>

$$N_{b_{x,y}}^{th} = \frac{32}{3\sqrt{2}\pi} \frac{R|\eta|\varepsilon_l^{2\sigma}\omega_r}{\langle\beta\rangle_{x,y} Z^2 \beta^2 c R_{\perp}}$$

if we increase the slip-factor, $|\eta|$, the bunch intensity limit would increase

- → Redesign of the DR lattice to increase $|\eta|$ which also increases the average beta function:

A. Chancé, next talk

$$\left. \begin{array}{l} \gamma_{tr} = 27 \\ \rightarrow \\ \gamma_{tr} = 18.57 \end{array} \right\} \left\{ \begin{array}{l} |\eta_1| = 0.00127 \rightarrow |\eta_2| = 0.00276 \\ \langle\beta\rangle_{y_1} = 173.64 \text{ m} \rightarrow \langle\beta\rangle_{y_2} = 160.4 \text{ m} \end{array} \right.$$

- → N_b^{th} increase by factor $(\eta_2/\eta_1)/(\beta_2/\beta_1)$

- Matching the bunch in the bucket:
 → Increase voltage by factor η_2/η_1

$$V_{rf} = 26.75 \text{ MV}$$

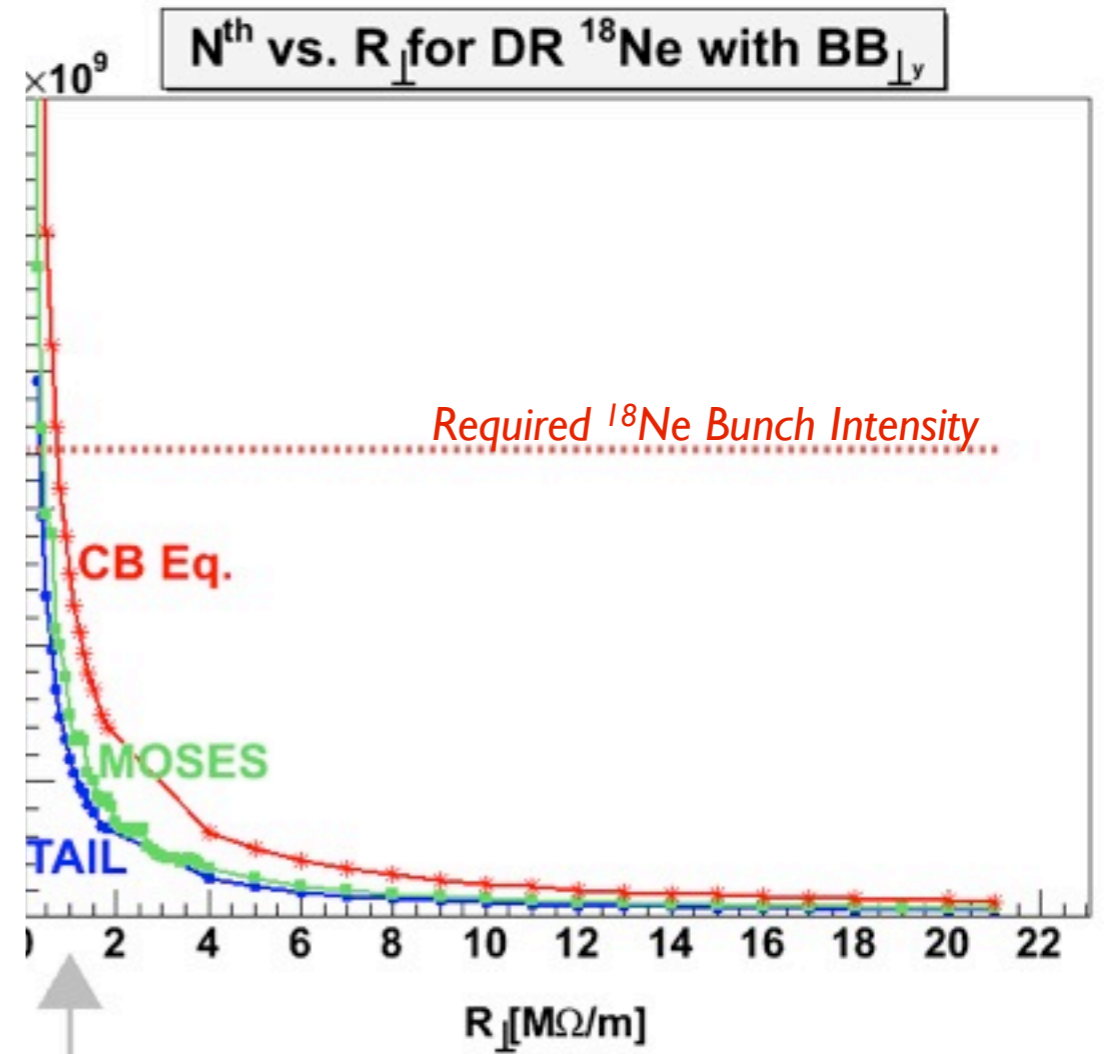
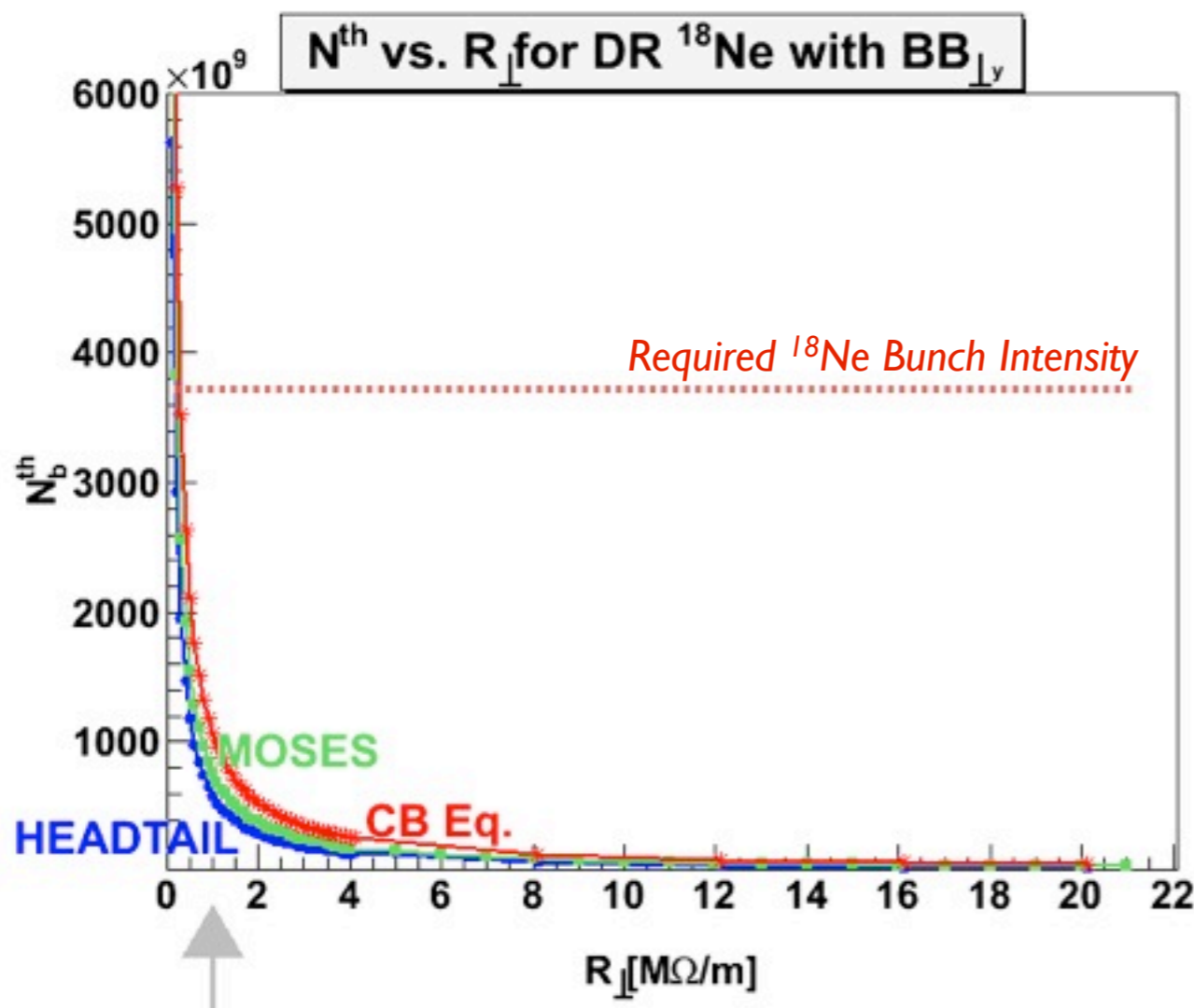
$$\frac{Q_s \beta c \tau_b}{2R|\eta|\delta_{max}} = 1$$

$$Q_s = \sqrt{\frac{hZeV_{rf}|\eta|}{2\pi\beta^2 E_{tot}}}$$

Decreasing γ_{tr} , Increasing V_{rf}

$$\begin{array}{lcl} \gamma_{tr} = 27.0 & \rightarrow & \gamma_{tr} = 18.6 \\ V_{rf} = 11.96 \text{ MV} & \rightarrow & V_{rf} = 26.75 \text{ MV} \\ L_{eff} = 36\% & \rightarrow & L_{eff} = 39\% \end{array}$$

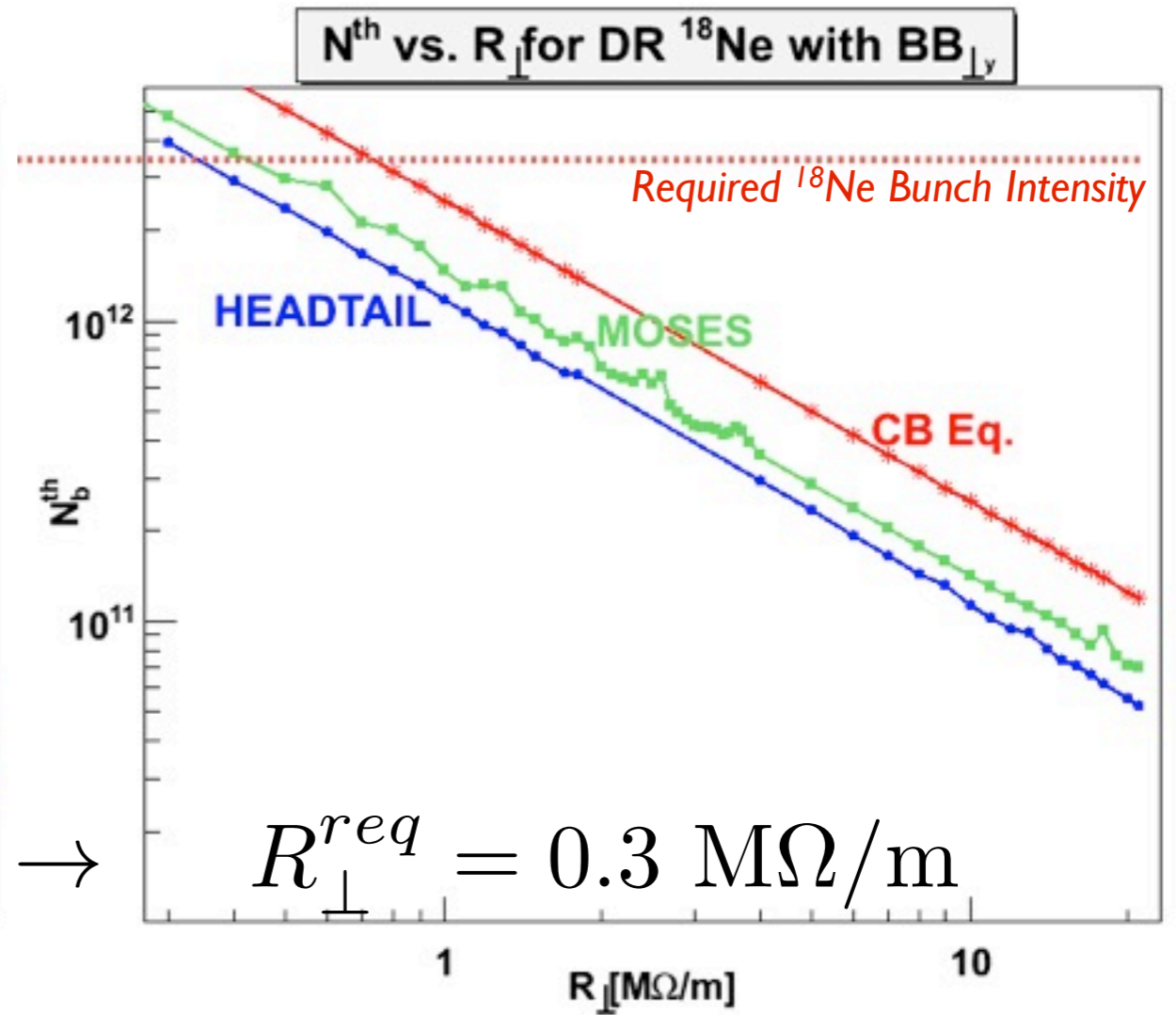
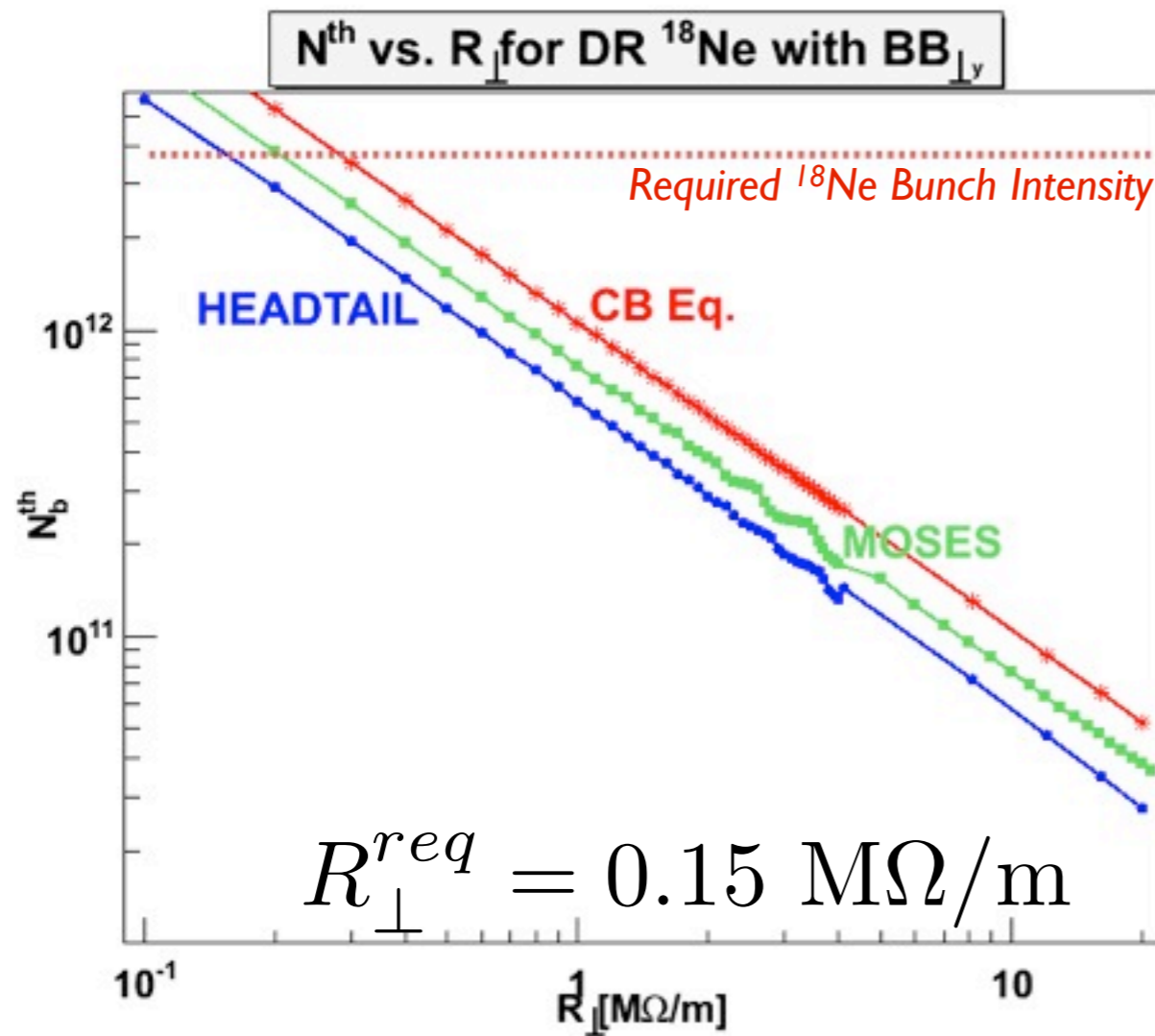
DRR REDSINGN



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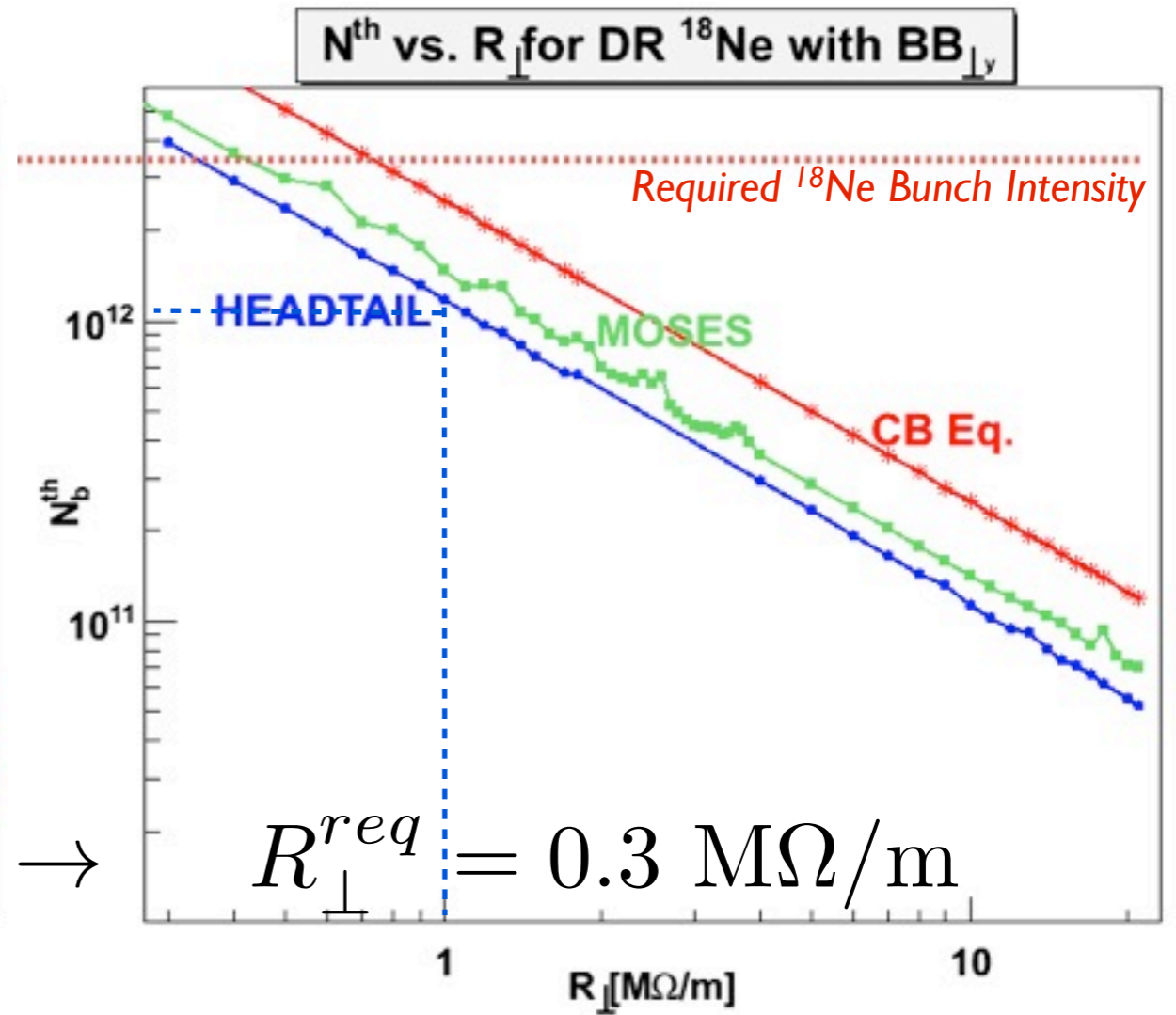
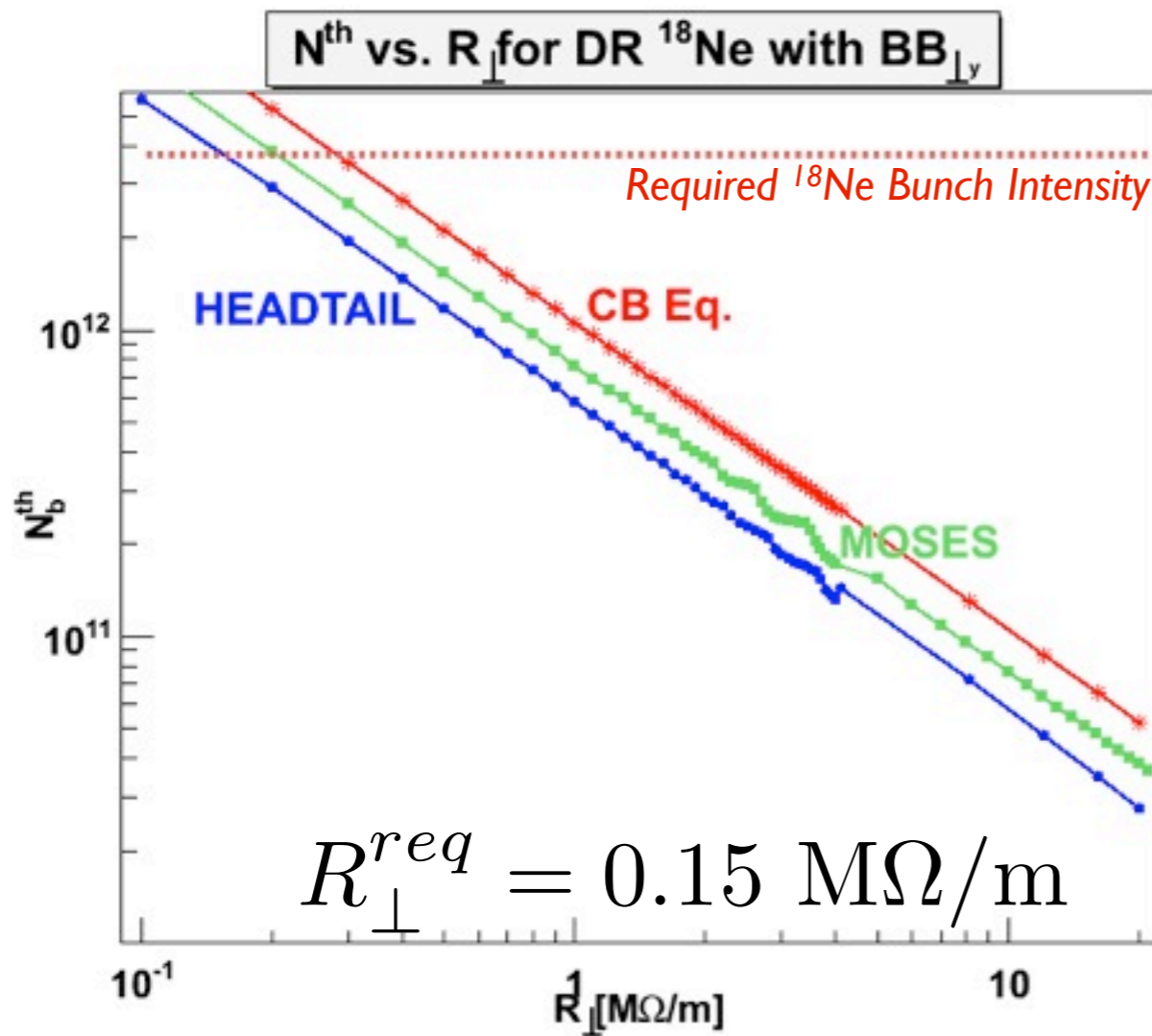
DRR REDSINGZ



- **So R_{\perp}^{req} is still a factor 3.3 too small ($R_{\perp}^{DR} \sim 1 \text{ M}\Omega/\text{m}$)**

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- $N_b^{th} = |2e|I|^{18}\text{Ne}$ can be used to get N_b^{th} for all other ions by using that **CB Eq.** goes as $N_{b_{x,y}}^{th} \propto \frac{A}{Z^2}$

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DR Intensity Limit for New Lattice

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$R_{\perp DR} =$ I MΩ/m	Bunch Intensity Limit, N_b^{th}	
	[e12]	[% of N_b^{nom}]
⁶He	10	224
¹⁸Ne	1.2	35
⁶He	10	112
¹⁸Ne	1.2	70
⁶He	10	112
¹⁸Ne	1.2	175
⁸Li	5.9	129
⁸B	2.1	127
⁸Li	5.9	65
⁸B	2.1	64
⁸Li	5.9	26
⁸B	2.1	25

Note; In Donini's table $SF = 10^{-4}$ while we are using $SF = 5 \cdot 10^{-3}$

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⁸ Li	⁸ B	⁶ He
129	127	224
65	64	112
26	25	45

$$N_b^{nom} = \frac{\Phi_0 L_{circ} t_{sps}}{N_{bunches} L_{eff} T_{eff}} \left(1 - 2^{-\frac{t_{sps}}{\gamma t_1/2}}\right)^{-1}$$

Conclusions

- Transversal Broad Band Impedance enforces redesign of the Beta Beam Decay Ring
(Other collective effects still to be studied)
- A new design of the Decay Ring (by A. Chancé)
 - ➔ Increases slip-factor, voltage and straight fraction
 - ➔ More of the Beta Beam scenarios are allowed
(assumed $R_{\perp}^{DR} = 1 \text{ M}\Omega/\text{m}$):

$R_{\perp}^{DR} =$ 1 M Ω/m	Bunch Intensity Limit, N_b^{th}	
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Thank You!

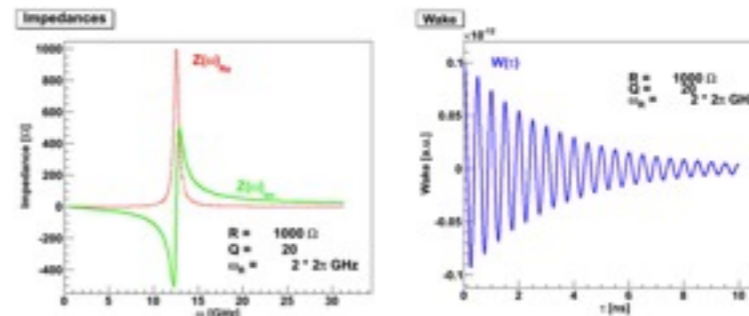


Backup Slides

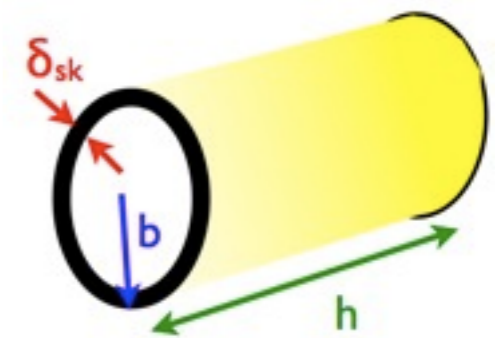
To Do

- Same study in longitudinal plane $Z_{||}(\omega) = \frac{R_{||}}{1 + iQ \left(\frac{\omega_r}{\omega} - \frac{\omega}{\omega_r} \right)}$
- ➔ Ongoing HEADTAIL simulations, but can't use MOSES since only for \perp

- Same with Narrow Band



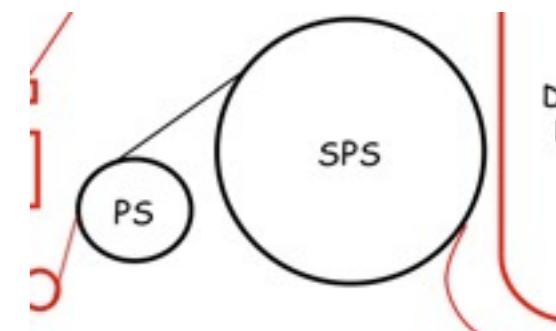
- Same with Resistive Wall Impedance



- Same with Space Charge



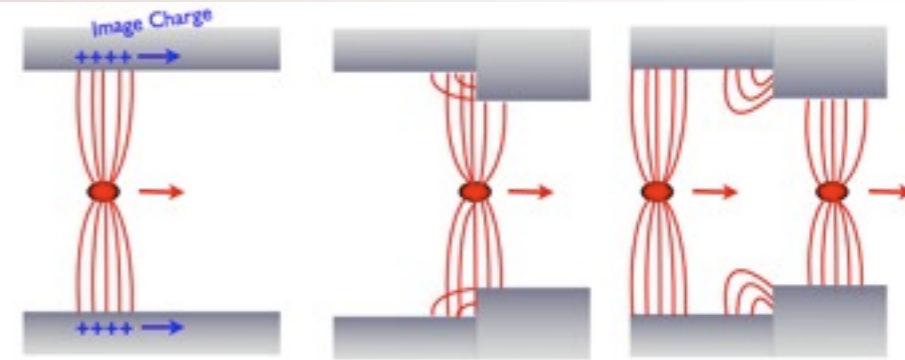
- Same with the already existing SPS & PS



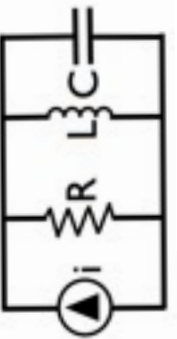
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Resonance Impedance

- Wake Fields (*time domain; $W(t)$*) can
 - be trapped in pipe cavities
 - cause “**Resonance Impedance**”



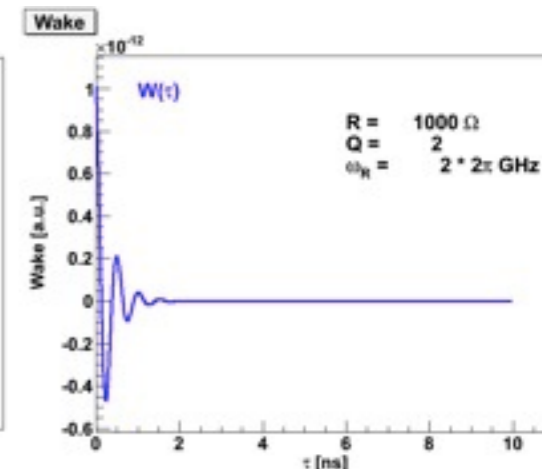
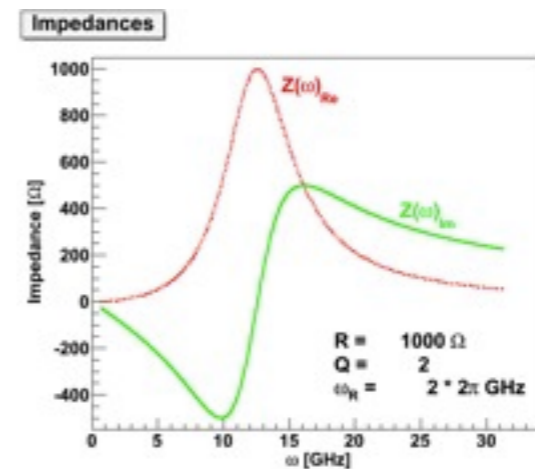
- Resonance Impedance (*frequency domain; $Z(\omega) = \mathcal{F}[W(t)]$*),
 - in the Transverse plane can be modeled by an RLC circuit as:



$$Z_{\perp}(\omega) = \frac{R_{\perp} \frac{\omega_r}{\omega}}{1 + iQ \left(\frac{\omega_r}{\omega} - \frac{\omega}{\omega_r} \right)}$$

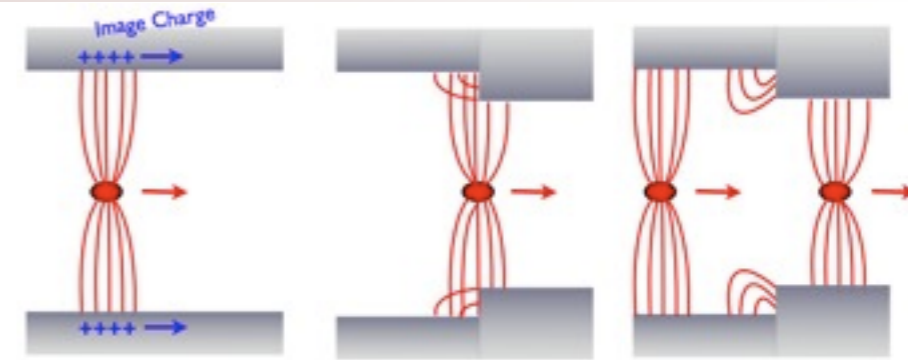
Q = “Quality Factor”
 ω_r = “Resonance Angular Frequency”
 R_{\perp} = “Shunt Impedance”

- For low Quality Factor ($Q \approx 1$) the Wake Field is short lived and the impedance is “Broad Band”

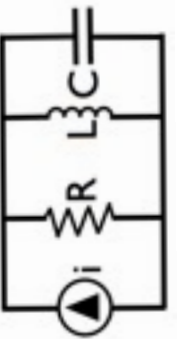


Resonance Impedance

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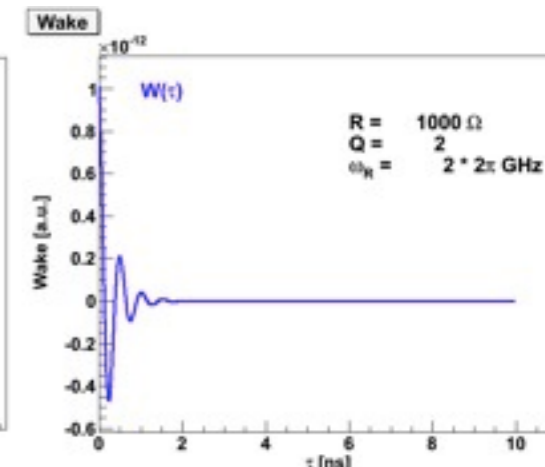
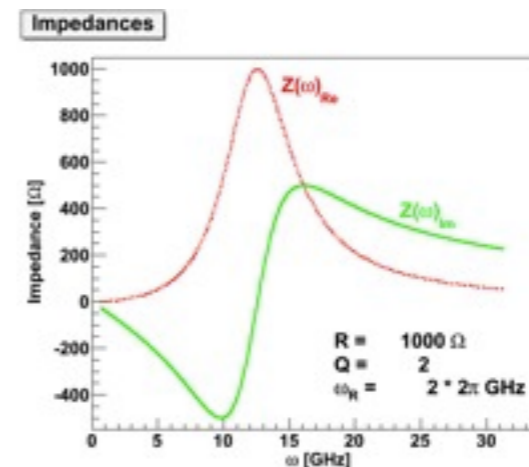
$$Z_{\perp}(\omega) = \frac{R_{\perp} \frac{\omega_r}{\omega}}{1 + iQ \left(\frac{\omega_r}{\omega} - \frac{\omega}{\omega_r} \right)}$$

Q = “Quality Factor”
 ω_r = “Resonance Angular Frequency”
 R_{\perp} = “Shunt Impedance”

= **I**
 $\approx \omega_c = \beta c / b_y$
 (see next slide)

- For low Quality Factor ($Q \approx 1$) the Wake Field is short lived and the impedance is “Broad Band”

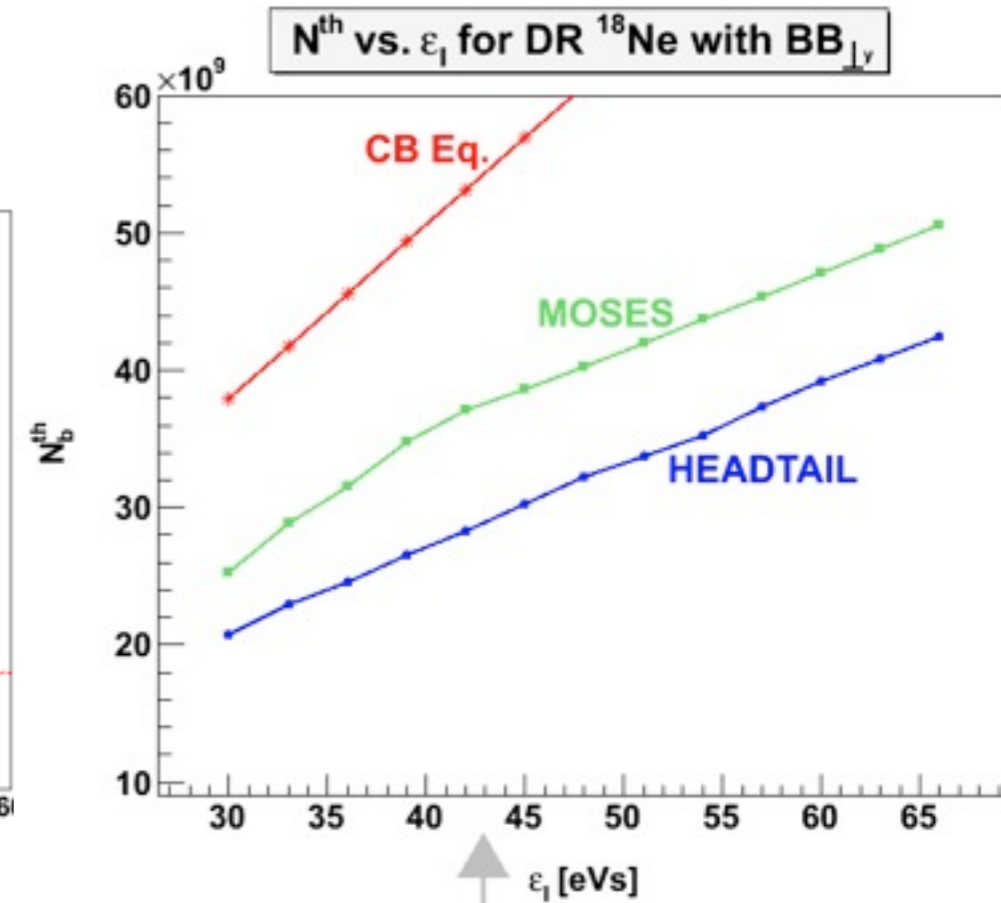
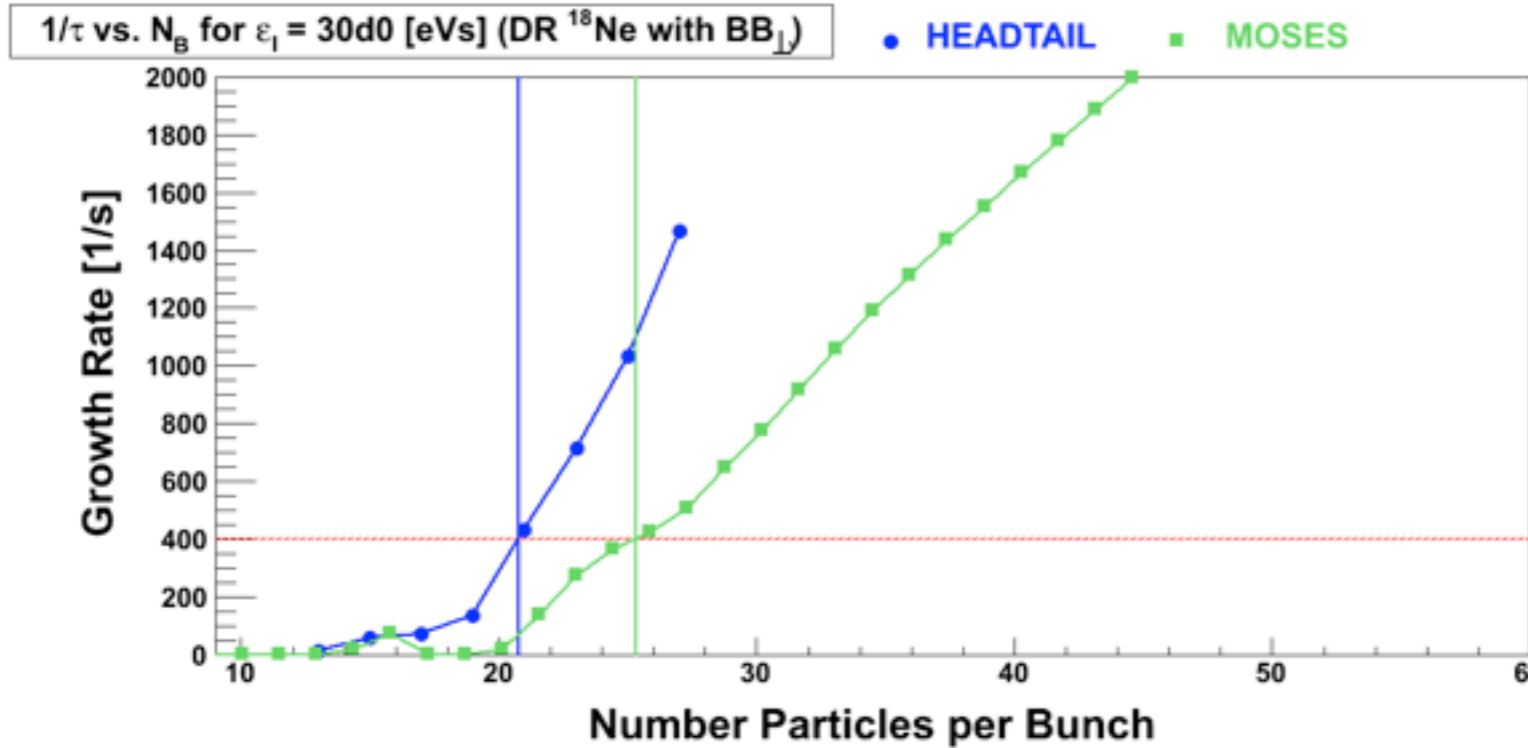
- Will show results from “**Transverse Resonance Broad Band Impedance**”



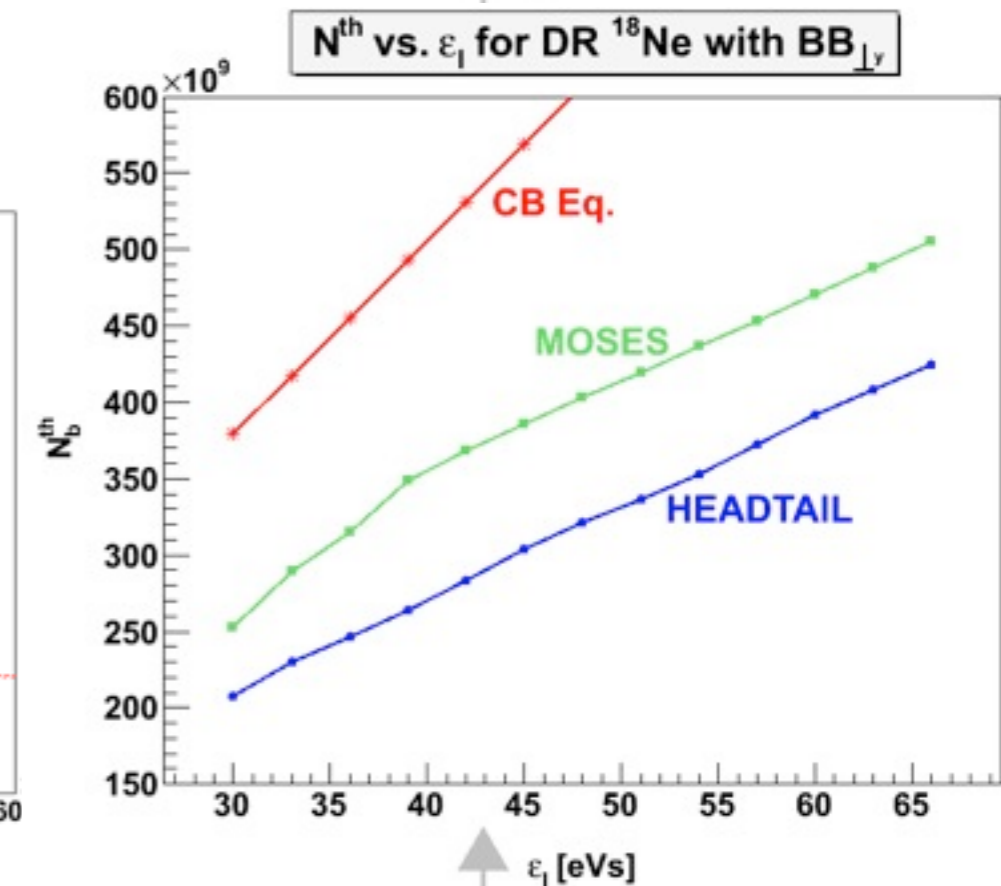
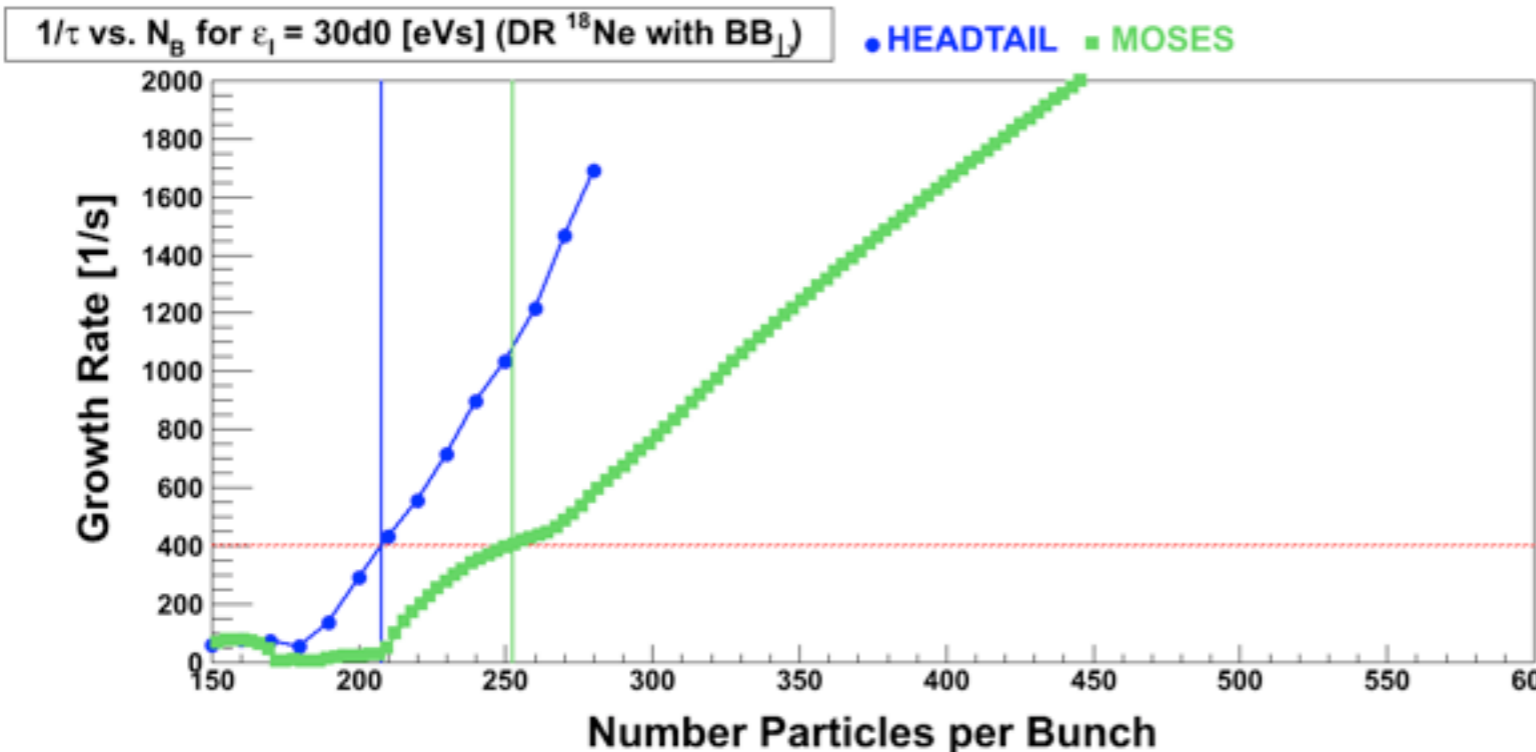
ϵ_1 Scan for ^{18}Ne

$(I/\tau)^{\text{th}} = 400\text{Hz}$

$R_{\perp} = 20\text{ M}\Omega/\text{m}$ (SPS)



$R_{\perp} = 2\text{ M}\Omega/\text{m}$ (RHIC)



T H R E S H O L D

Shunt Impedance

- The “Shunt Impedance”, R_{\perp} , is the main parameter in the RLC model of the Resonance Impedance

$$Z_{\perp}(\omega) = \frac{R_{\perp} \frac{\omega_r}{\omega}}{1 + iQ \left(\frac{\omega_r}{\omega} - \frac{\omega}{\omega_r} \right)}$$

- Modeling existing machines the same way we have

	PS	SPS	LHC (at top energy)	LHC (no collimators)	RHIC
R_{\perp} [MΩ/m]	3	20	30	2	2

E. Métral

W. Fisher et. ali, Analysis of Intensity Instability Threshold at Transition in RHIC