

It was remarked that the spatial luminosity profiles $L(z)$ as measured by the experiments are looking very much Gaussian. From this fact it might be concluded that bunches in the machine *must* also have Gaussian profile. We will show that this conclusion is not correct for realistic data and that – even under the assumption of perfectly identical bunches – a de-convolution of $L(z)$ will only give very imprecise bunch profiles.

It is well known that for two identical Gaussian bunches $L(z)$ will be precisely Gaussian, then with half the σ^2 of the bunches. However, we will now calculate $L(z)$ for a \cos^2 bunch profile that is often used to describe long bunches in proton machines¹ and that significantly deviates from a Gaussian shape with the same ‘nominal bunch length’, especially in the tails, Fig. 1.

The \cos^2 profile can be described by the parameter x with $-0.5 \leq x \leq 0.5$ relative to the full bunch length

$$\rho(x) \propto \begin{cases} \cos^2(x \cdot \pi) & [-\frac{1}{2} \leq x \leq +\frac{1}{2}] \\ 0 & [else] \end{cases} = \begin{cases} (1 - \cos(x \cdot 2\pi))/2 \\ 0 \end{cases}$$

For simplification we have not normalized $\rho(x)$, not important in the present context.

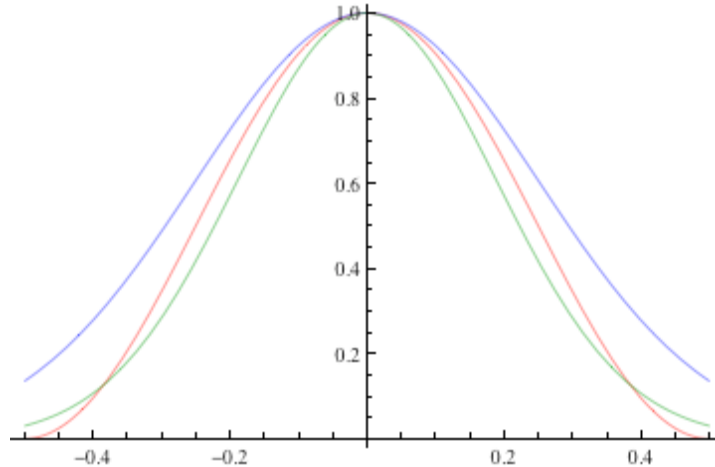


Fig.1: \cos^2 bunch profile (red), Gaussian profile with the same nominal “ 4σ bunch length” (blue) and Gaussian profile reproducing best the \cos^2 luminosity profile (green), see later.

The luminous profile is the density product of both bunches at a given z integrated over time. We assume two identical bunches both moving with $v \approx c$ in opposite direction. Without loss of generality the variables t and z are defined such that bunch centers meet at $z=t=0$ and the absolute bunch length is normalized to 1, i.e.

$$L(z) \propto \int_{-\infty}^{+\infty} \rho(z - c \cdot t) \cdot \rho(z + c \cdot t) dt$$

Taking the above \cos^2 profiles and considering that ρ is zero outside the given x -range, $L(z)$ becomes here proportional to

¹ in contrast to the short Gaussian bunches in synchrotron radiation dominated electron machines

$$L(z) \propto \int_{(-1/2+|z|)/c}^{(1/2-|z|)/c} (1 + \cos(2\pi(z - c \cdot t))) \cdot (1 + \cos(2\pi(z + c \cdot t))) dt$$

which is mathematically equal to

$$L(z) \propto \frac{(2 \cdot |z| - 1)(2 + \cos(4\pi |z|)) + 3 \sin(4\pi |z|)}{4\pi}$$

We did not do a mathematical fit for the best matching Gauss curve for this $L(z)$ but simply played with numbers and plotting the results. The closest match was obtained for $\exp(-28 \cdot z^2)$, see Fig. 2, corresponding to $\sigma_L = 1/\sqrt{56} \approx 0.134$. One sees (Fig. 2) that both curves are nearly indistinguishable, hence it is nearly impossible from real data to de-convolute the true bunch profile from $L(z)$. This Gaussian luminosity profile would be created perfectly by Gaussian bunches of $\sigma_B = 1/\sqrt{28} \approx 0.19$, i.e. bunches would have a “ 4σ bunch length” of 0.75 compared to a bunch length of 1 for the \cos^2 bunches.

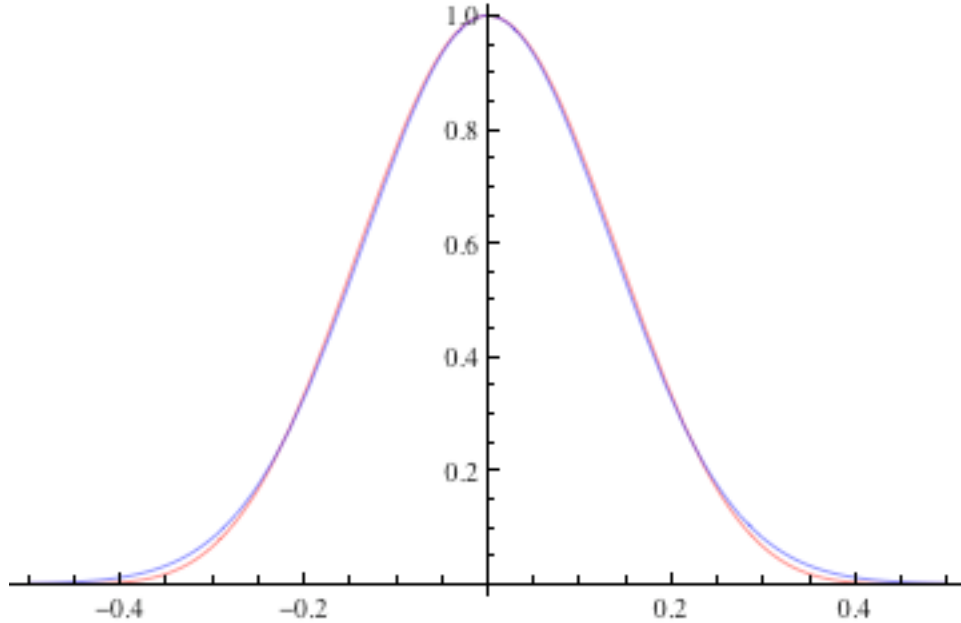


Fig.2: $L(z)$ for \cos^2 -bunches (red, $4\sigma=1$) and Gaussian bunches (blue). The Gaussian profile drawn is $\exp(-28 \cdot z^2)$.

On top of these facts, established for two absolutely identical bunches, one should consider that there is a certain bunch-to-bunch scatter in shape which in general makes things look ‘even more Gaussian’. This is true for any two interacting bunches, if only one bunch-pair would be gated out, and even more for the integral $L(z)$ measured over the whole set of bunches (full beam).

Conclusion: It is practically impossible to de-convolute the true bunch-profile from the longitudinal luminosity profile even if measured very precisely by the experiments. This statement even holds when the measurement would be done in gating only a single bunch pair and assuming that both bunches have precisely the same longitudinal density distribution.

Comparison with LHC measurements:

Form Ph. B.: The experiments measure – integrated over all bunch crossings – a Gaussian luminous region with $\sigma_L=60$ mm (\pm a few mm for the different experiments, but very stable along the coast for the bunches now with ‘calibrated length’ by blow-up during the ramp), i.e. $\sigma_L=180$ ps in ‘time domain’. From this one would conclude that bunches – if assumed Gaussian – would have $\sigma_B=180\cdot\sqrt{2}$ ps = 255 ps to create such a luminous region. On the other hand the RF claims from the ‘Juliameter’ a 4σ bunch length² of 1350 ps or $\sigma=338$ ps. The ratio of these is 0.75, exactly the ratio of the \cos^2 full bunch length and the best-fitting (luminosity wise) Gaussian type bunches as obtained above.

This rises (again) the discussion about what is “the” bunch length (and, maybe, if a single number is sufficient to grasp the essential quantities of a bunch)

² which is in fact measured as full-width-at-half-maximum and then rescaled to the nominal 4σ bunch length