Electron Cloud in IRs?

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Main Factors

- E-cloud influences incoherent and coherent oscillations of beam particles in various aspects.
 - It works as a static lens, shifting up all coherent and incoherent tunes.
 - It gives a significant tune spread. With the size of the e-cloud similar to the proton beam size, the nonlinear tune spread is comparable to the tune shift. The tune spread is defocusing with the amplitude.
 - As a reactive medium, e-cloud works as a sort of low-Q impedance at the electron bounce frequency ω_e which phase advance on the bunch rms length is $\omega_e \approx 0.5 \sqrt{N \pi \sigma / \sigma^2}$

$$\psi_e \cong 0.5 \sqrt{N_b r_e \sigma_z} / \sigma_x^2$$

• Note that number of e-cloud pinches per p-bunch is $\sim \psi_e$. Thus, for $\psi_e \gtrsim 1$ the effective size of the electrons within the proton beam is ~2-3 times smaller than the proton bunch radius.

Assuming e-cloud transverse profile same as for the beam, the incoherent tune follows:



Wake function

 Following [Burov & Dikansky, 1997], e-cloud wake can be modeled as a low-Q resonator:

$$W(\tau) \simeq W_0 \sin(\omega_e \tau) \exp(\omega_e \tau / 2Q);$$
$$W_0 = \frac{2N_e r_e c}{\sigma_\perp^4 \omega_e}, \quad Q \sim 2 - 3, \ \tau < 0$$

equivalent to a shunt impedance

$$\frac{R_s}{Q} \simeq Z_0 \frac{N_e r_e c}{2\pi \sigma_\perp^4 \omega_e}; \ Z_0 = \frac{4\pi}{c} = 377 \text{ Ohm}$$

Here $\frac{N_e}{r}$ is number of electrons seen inside the proton beam size of the radius σ_{\perp} per revolution.

Weak Head-Tail (WHT)

 Application this wake function to the WHT tune shift and growth rate (A. Chao, Eq. 6.213, air-bag) results in (HT phase χ≤1):

$$\operatorname{Im}[\Delta Q] = \chi \Delta Q_{e0} F_R(m, \phi_e); \quad \operatorname{Re}[\Delta Q] = \Delta Q_{e0} F_I(m, \phi_e); \quad \phi_e = \sqrt{2\omega_e \sigma_z} / c$$

Growth rates factors vs BB wake phase advance

Mode tune shifts factors vs BB wake phase advance



$$F_{R}(m,Q,\phi) = 3\int_{0}^{\infty} \frac{J_{m}(x\phi)J_{m}'(x\phi)}{1+Q^{2}(x-1/x)^{2}} \frac{dx}{x}; \quad F_{I}(m,Q,\phi) = \frac{3}{2}\int_{0}^{\infty} \frac{Q(x-1/x)J_{m}^{2}(\phi x)}{1+Q^{2}(x-1/x)^{2}} \frac{dx}{x}$$

Unstable modes have positive tune shift – thus, they are not L-damped after the SD shift s to the left due to e-cloud unharmonicity!

At $\chi \simeq 1$, for the MUM: $\text{Im}[\Delta Q] \approx \text{Re}[\Delta Q] \approx 0.2 \Delta Q_{e0}$ (used for markers at p.3 plot)

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Why End of the Squeeze?

According to the air-bag plots of p.5, the most unstable mode (MUM) number *l* ≅ 1.4*\nu*_e. To be not suppressed from the longitudinal L-damping this number cannot be too high, *\nu*_e ≤ 1-2. For the growth rate do not be too low, the electron phase advance cannot be too small as well, *\nu*_e ≥1. Thus, to drive the instability, the phase has to be about 1:

 $\Psi_e \sim 1$



BB, LO=500A, dQe0=8.0E-4

 $N_e = 1.3 \cdot 10^{10}$ total

- During the squeeze, the phase advance Ψ_e significantly changes: $\psi_e = \begin{cases} 9 \text{ rad for } \beta = 300 \text{m} \\ 2 \text{ rad for } \beta = 4 \text{km} \end{cases}$
- Thus, the effective number of electrons $\frac{N_e}{N_e}$ has to be shared between 4 high –beta regions of IR1 and IR5, requiring ~4E9 e per every ~25m region.
- Due to the e-pinches, this number has to be expected 2-10 times lower.

SD Focusing-Collapse for LO=0



LO=+140A – computed threshold

BB only

BB, LO=0, dQe0=2.4E-4

BB, LO=0, dQe0=4.8E-4

BB, LO=0, dQe0=7.2E-4

For zeroed LO, it takes twice less electrons for the instability than for +500A.

SD Collapses and Reductions for LO=-500A, No EC



No D-collapse, just a minor reduction... This may be enough to drive the instability at Jura - but not at the Plateau.

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<u>SD Collapses and Reductions for LO=-500A</u>



- The instability may develop at adjust only not at the squeeze.
- With that polarity, the beam is stable both at too low and too high e-cloud.
- Instead, for LO>0, there is no stabilization with increase of e-cloud.

Arguments for IR e-cloud hypothesis

- IR e-cloud gives an instability mechanism, sensitive to 2 beams and not requiring coupled-beam oscillations.
- This instability appears only at the end of the squeeze, which is consistent with decrease of the phase advance ψ_e during the squeeze:

 $\psi_e = \begin{cases} 9 \text{ rad for } \beta = 300 \text{m} \\ 2 \text{ rad for } \beta = 4 \text{km} \end{cases}$

- It takes only a few E9 e/IR to make the instability possible.
- Coupled-beam is refuted both conceptually and experimentally, and we do not know any other 2-beam instability mechanism.
- It is reasonable to expect only one beam oscillating (which was observed).
- It is reasonable to expect both emittances degrading (observed at cogging MD)

10

This hypothesis is consistent with LO<0 observations (S. Fartoukh)

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Questions

- What assumptions have to be taken for the IR vacuum chamber to make possible accumulation of a few E9 e/IR with 2 beams there?
- Can this Ne/IR be consistent with our knowledge about the IR?

 In case of no-refutation from the build-up simulations, can we install anti-cloud solenoids outside the IR quads?



